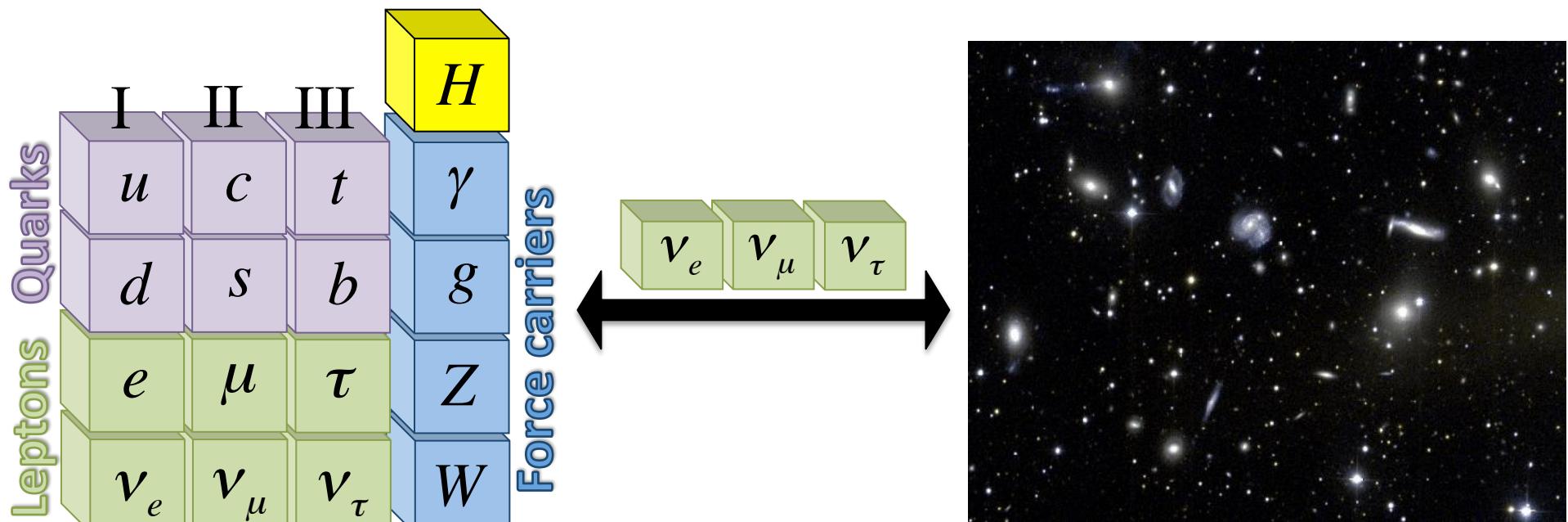


Massive neutrino signatures on the large-scale structure of the Universe

Francisco Villaescusa-Navarro
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Outline

1. Introduction
2. Cosmic neutrino background
 - Effects at linear order
 - Effects at non-linear order
3. Conclusions

Neutrinos

Pauli 1930

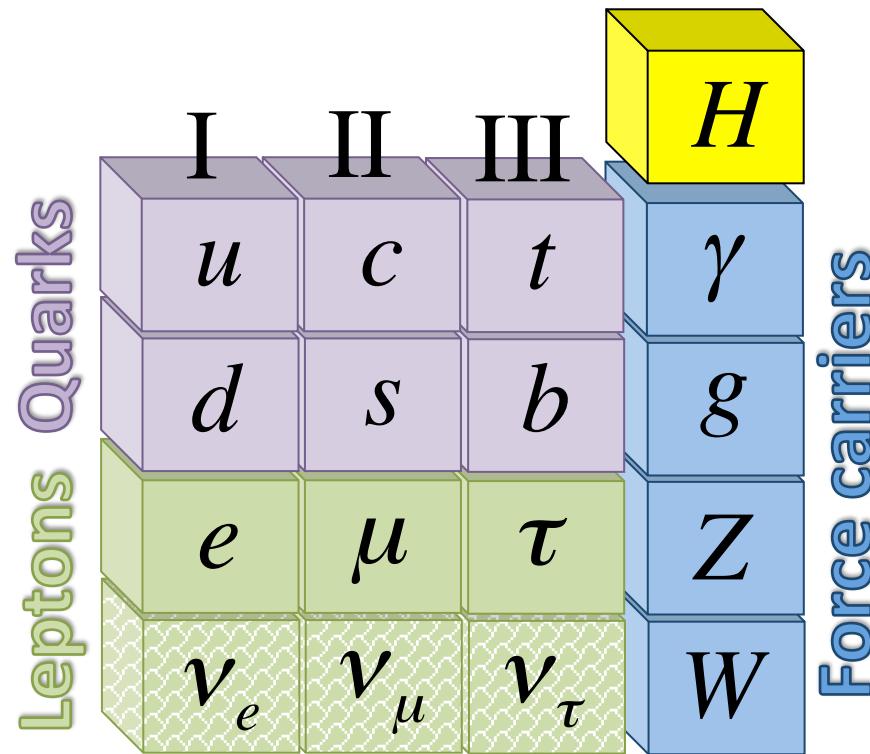
$$n \rightarrow p^+ + e^-$$

Violation of Energy
Violation of Momentum
Violation of Spin

$$n \rightarrow p^+ + e^- + \bar{\nu}_e$$

Cowan & Reines 1956

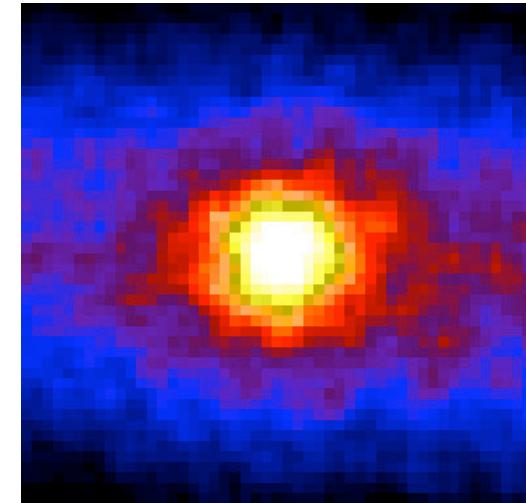
- Fundamental particles
- Neutral leptons: $\frac{1}{2}$ spin fermions
only sensitive to electroweak force and *gravity*
- $N_\nu=3$ from Z boson decay
- Very weak cross section
a neutrino could pass through a light year of lead
and not be stopped by any of the lead atoms!
- Massless in the SM



Neutrinos

Implications

- Physics beyond the standard model
- Cosmology
(neutrinos are the second most abundance particle in the Universe)
- Neutrino astronomy
(account for neutrino oscillations to estimate expected fluxes)



Open questions

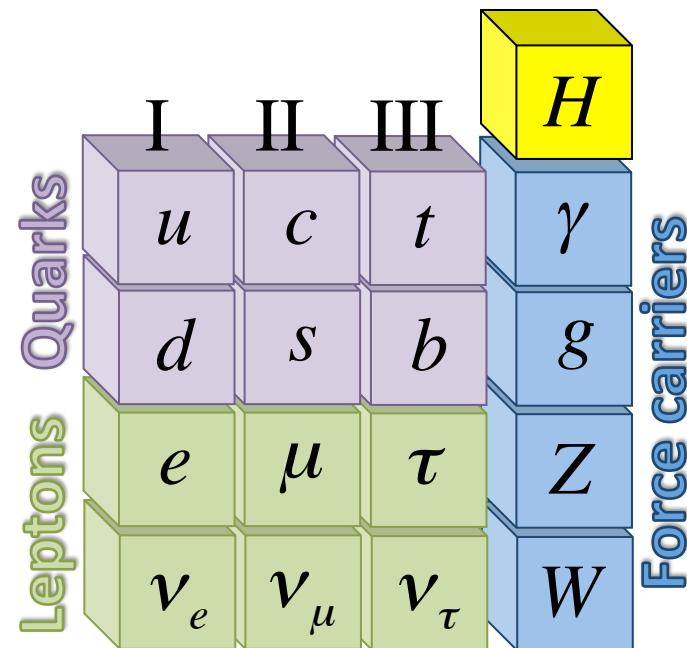
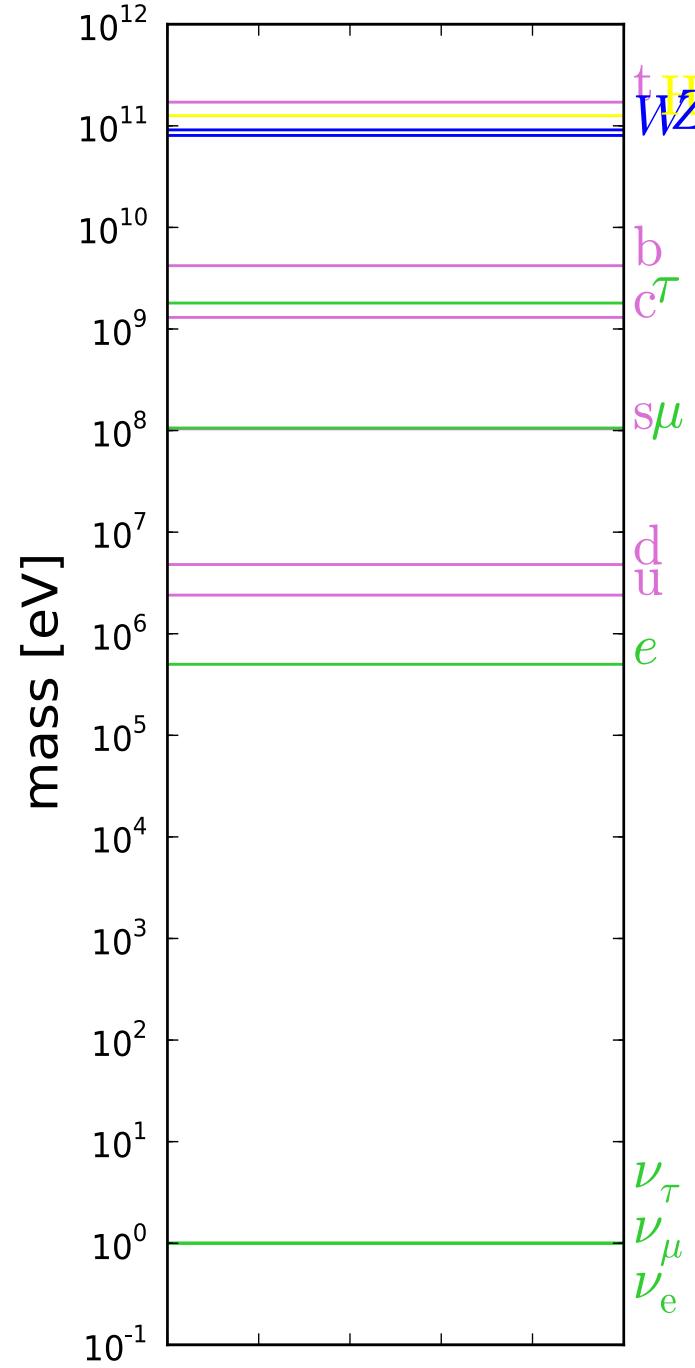
- What is the neutrino mass hierarchy?
- What is the absolute neutrino mass scale?
- What is the neutrino nature? Dirac or Majorana?

Neutrinos

- Neutrinos have mass
- What are the neutrino masses?

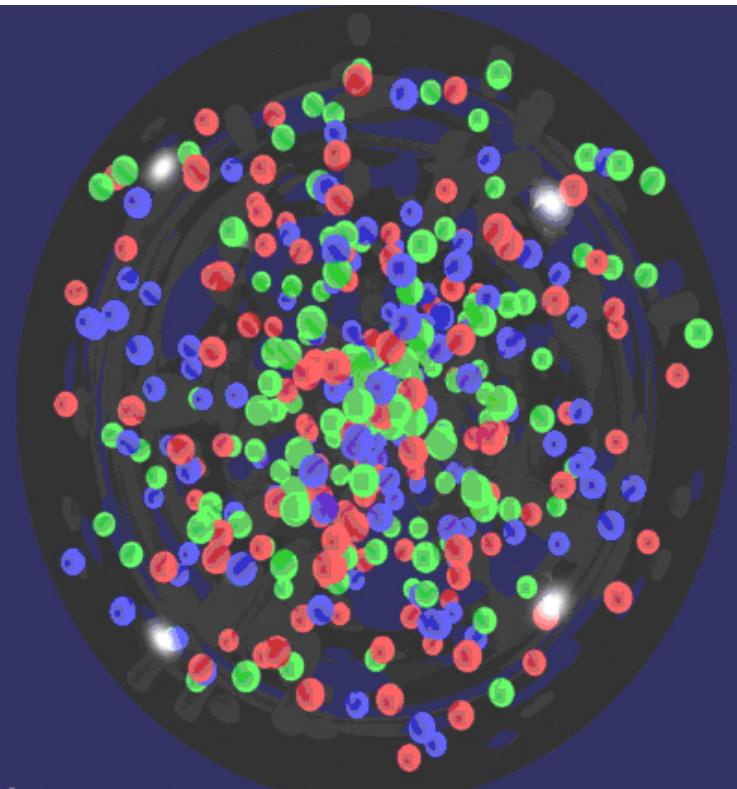
$$\sum m_\nu \geq 0.06 \text{ eV} \quad m(\bar{\nu}_e) \leq 2.3 \text{ eV}$$

1. *Impact on cosmology*
2. *Weigh neutrinos with cosmology*



Cosmic neutrino background

$$T \approx 10^{10} \text{ K} \approx 1 \text{ MeV}$$



- Photons
- Electrons & Positrons
- Neutrinos & Antineutrinos

$$\nu + e^- \Leftrightarrow \nu + e^- \quad e^- + e^+ \Leftrightarrow \nu + \bar{\nu}$$

$$n_\nu(p, z) dp \cong \frac{4\pi g_\nu}{(2\pi\hbar c)^3} \left(\frac{p^2 dp}{e^{(p/k_B T_\nu(z))} + 1} \right)$$

$$T_{\nu,0} = \left(\frac{4}{11} \right)^{1/3} T_{\gamma,0} \approx 1.95 \text{ K}$$

$$T_\nu(z) = T_{\nu,0}(1+z)$$

Cosmic neutrino background

Properties

$$\bar{n}_\nu(z) = \int_0^\infty n_\nu(p, z) dp \cong 113 (1+z)^3 \frac{\nu}{cm^3}$$

$$\bar{V}_\nu(z) = \frac{1}{n_\nu(z)} \frac{1}{m_\nu} \int_0^\infty n_\nu(p, z) p dp \cong 160(1+z) \left(\frac{eV}{m_\nu} \right) \text{ km/s}$$

$$\rho_\nu = \int_0^\infty n_\nu(p, z) dp \sqrt{p^2 + m_\nu^2} \quad P_\nu = \int_0^\infty n_\nu(p, z) dp \frac{p^2}{3\sqrt{p^2 + m_\nu^2}}$$

At high redshift

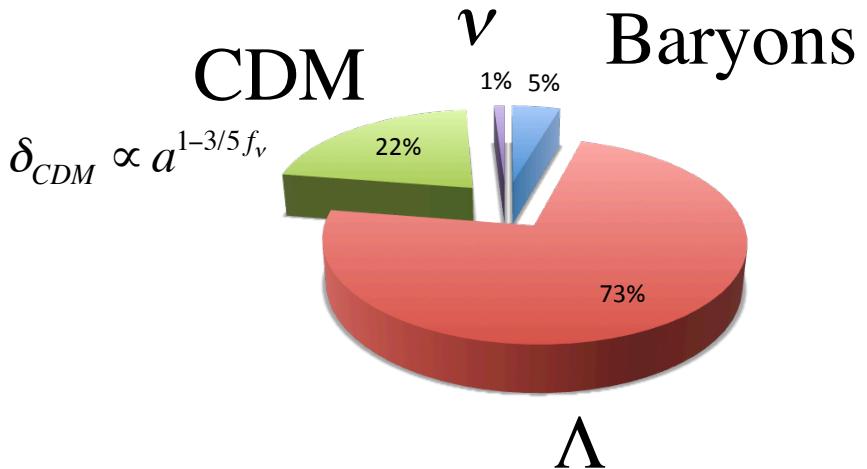
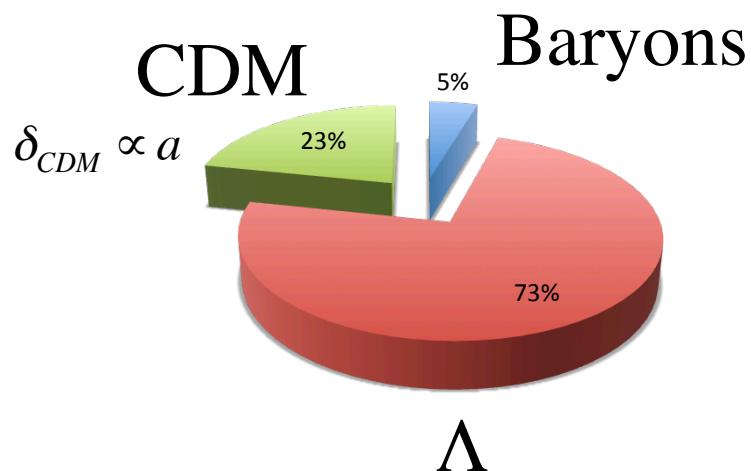
$$P_\nu \approx \frac{\rho_\nu}{3}$$

At low redshift

$$P_\nu \approx 0$$
$$\Omega_\nu h^2 = \frac{\sum_i m_{\nu_i}}{93.3 eV}$$

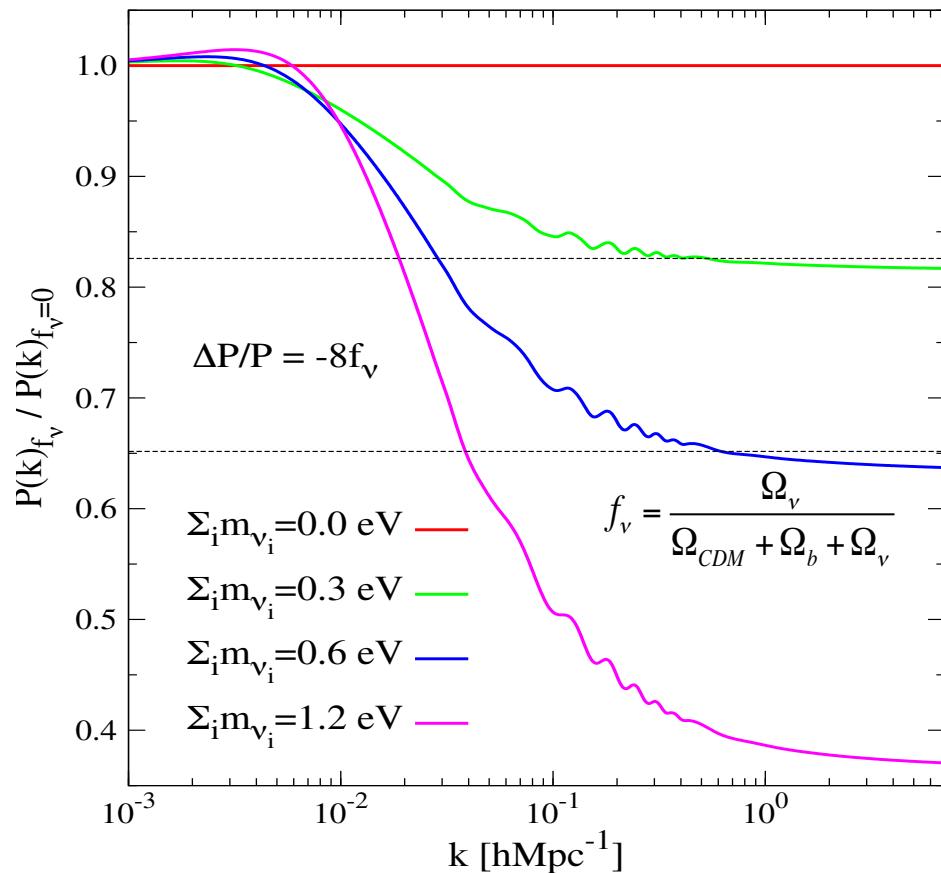
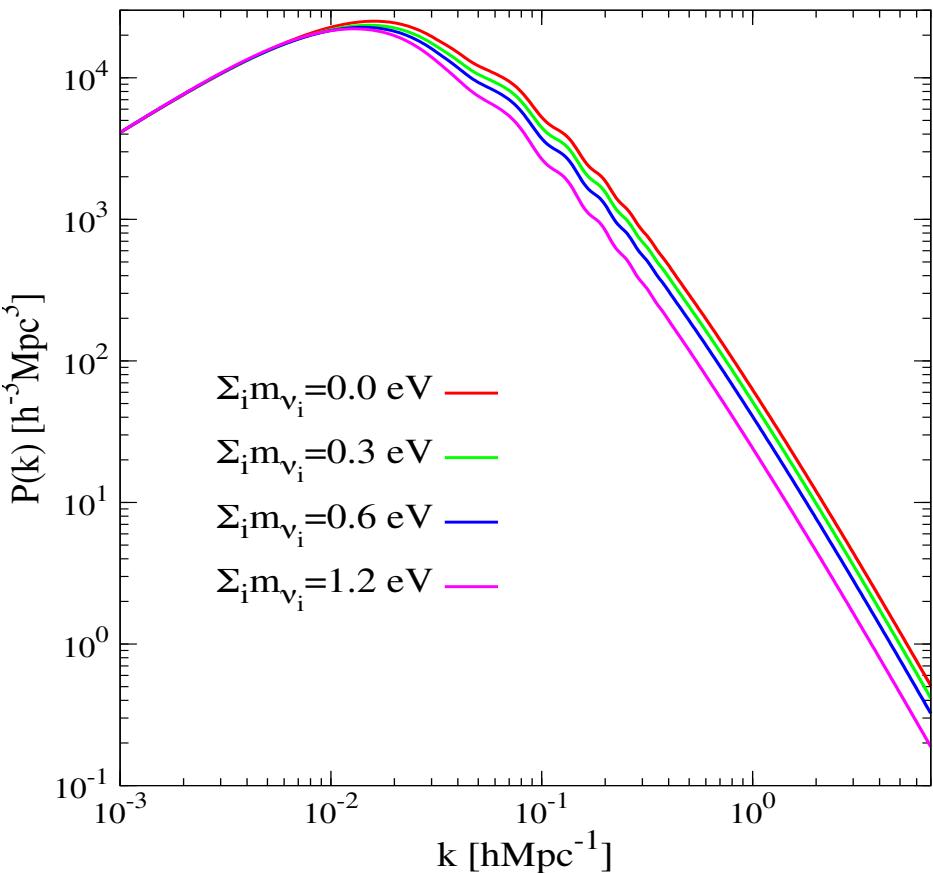
Effects at linear order

1. Modification of the Matter-Radiation equality time
2. Slow down the growth of matter perturbations



$$\Omega_{\text{DM}} = \Omega_{\text{CDM}} + \Omega_\nu = \text{fixed}$$

Effects at linear order

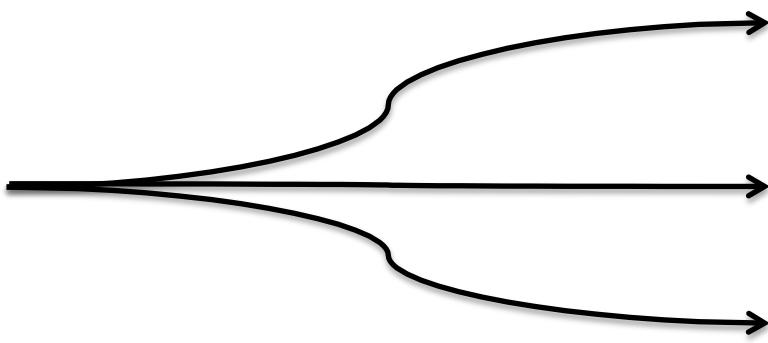


1. Modification of the Matter-Radiation equality time
2. Slow down the growth of matter perturbations

$$\Omega_{\text{DM}} = \Omega_{\text{CDM}} + \Omega_\nu = \text{fixed}$$

Effects on the non-linear regime

Why?

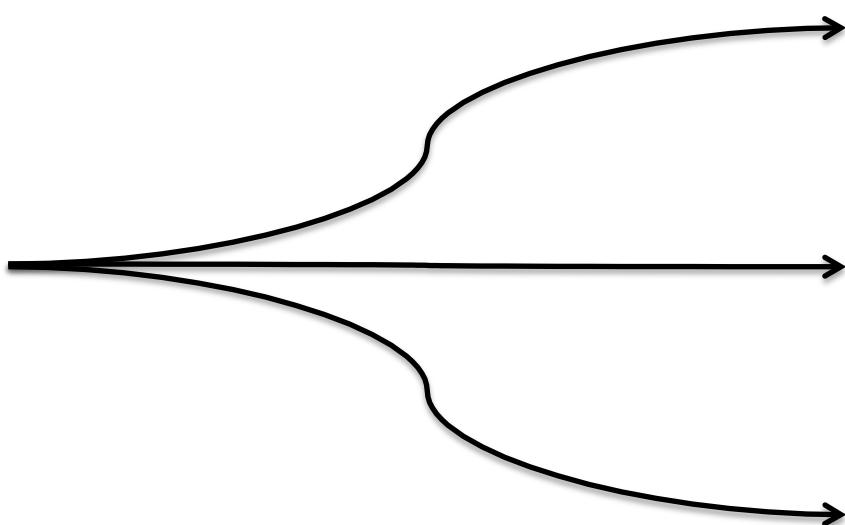


Important on small scales

Important at low redshift

Lots of modes in the
middle—non-linear regime

How?



Semi-analytic methods

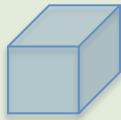
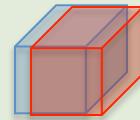
N-1-body, Ringwald & Wong, 2004
(see also Singh & Ma, 2003)

Perturbation theory

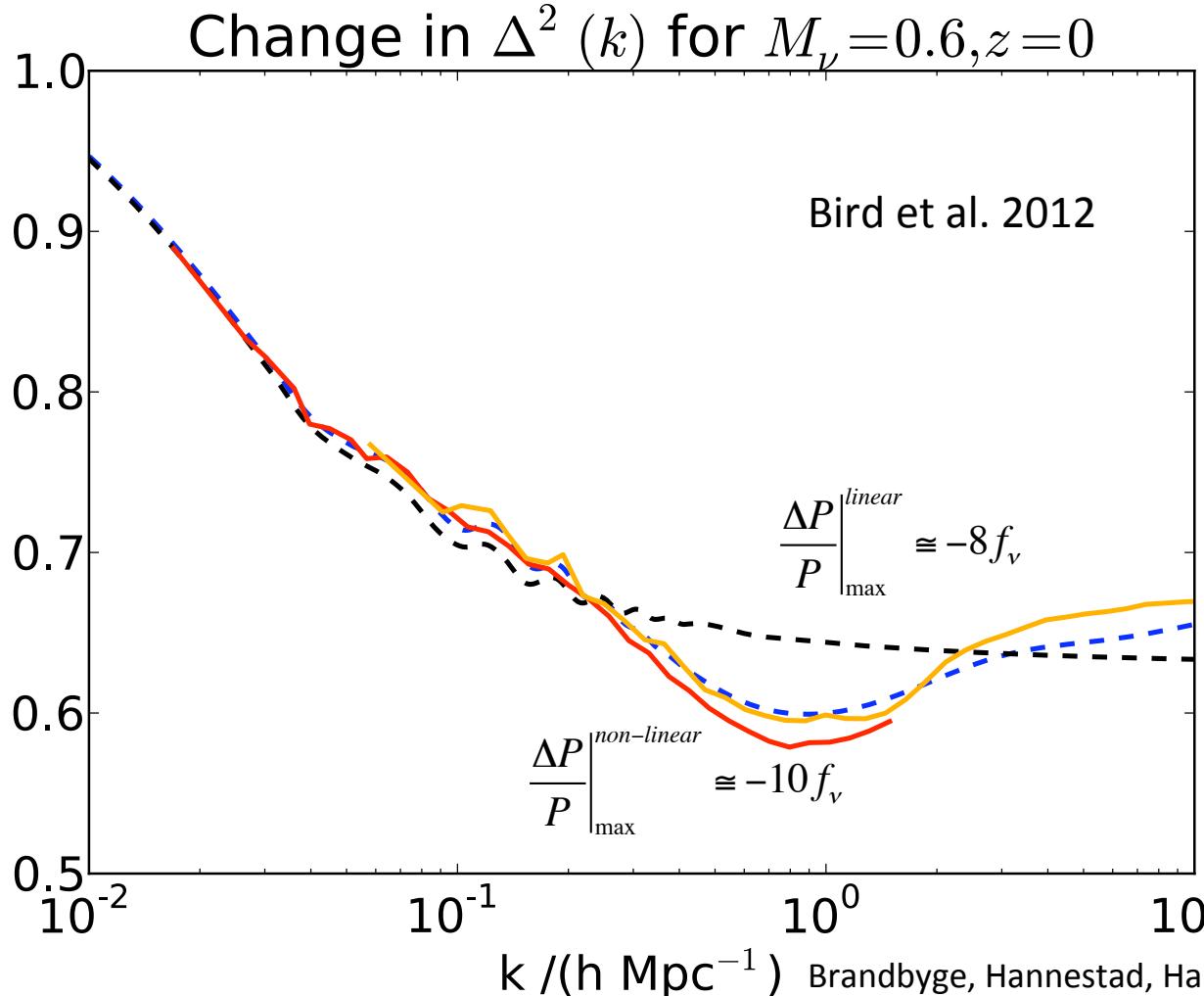
Blas, Garny, Konstandin, Lesgourges, 2014

N-body simulations

N-body simulations with neutrinos

	CDM	CDM + ν
<u>Power spectrum</u>	$P_m(k)$ 	$P_{cb}(k)$ $P_\nu(k)$ 
<u>Growth factor</u>	Scale independent	Scale dependent
<u>Growth rate</u>	Scale independent	Scale dependent
<u>Velocities</u>	Peculiar	Peculiar Peculiar + thermal
<u>Radiation</u>	-	May be important

Effects on matter power spectrum

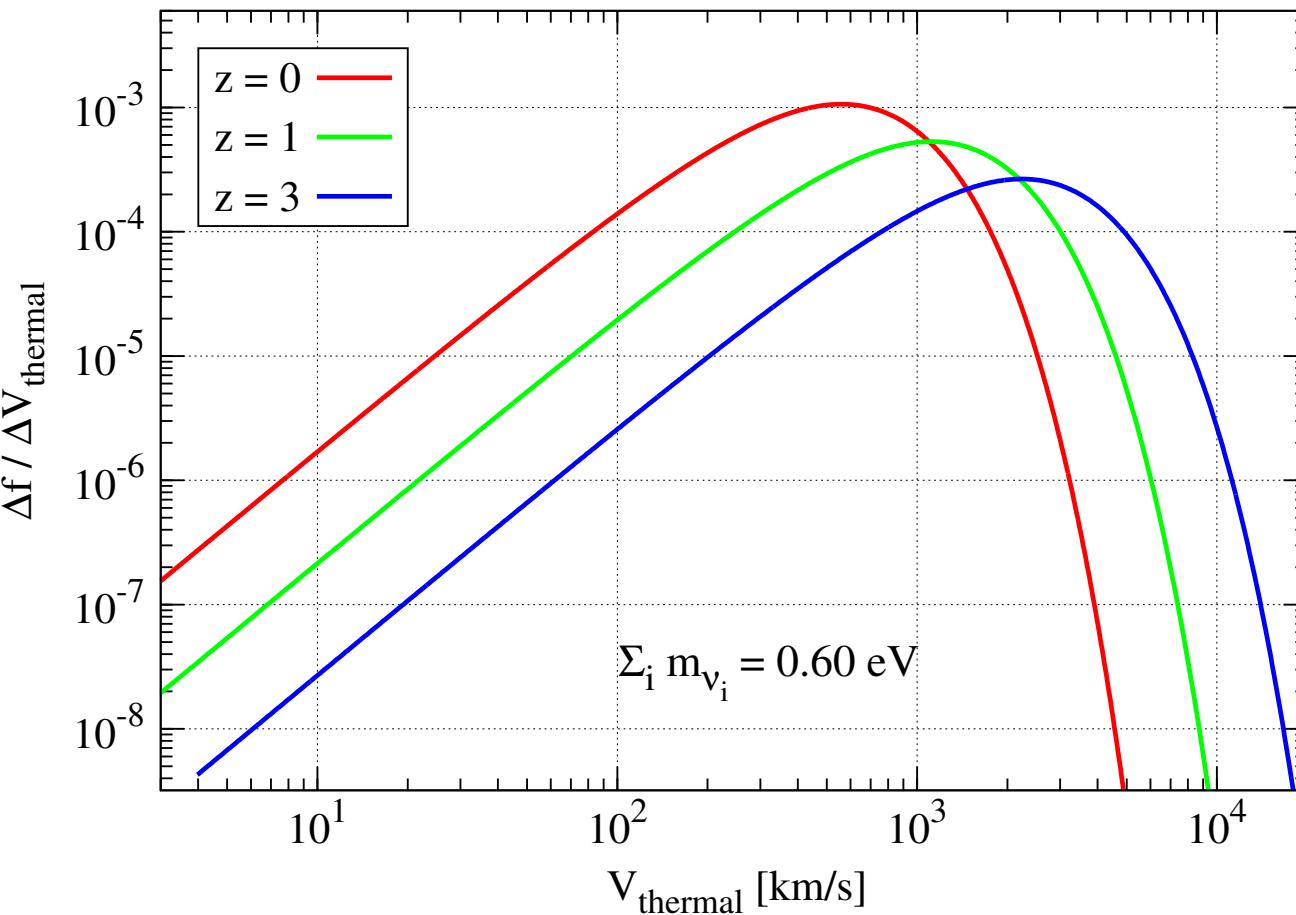


- Brandbyge, Hannestad, Haugbolle, Thomsen, 2008
Viel, Haehnelt, Springel, 2010
Bird, Viel, Haehnelt, 2012
Wagner, Verde, Jimenez, 2012
Agarwal, Feldman, 2011
Massara, FVN, Viel, 2014
Upadhye, Biswas, Pope, Heitmann, Habib, Finkel, Frontiere, 2014
Inman, Emberson, Pen, Farchi, Yu, Harnois-Deraps, 2015

Neutrino clustering

$$n_\nu(p, z) dp \cong \frac{4\pi g_\nu}{(2\pi\hbar c)^3} \left(\frac{p^2 dp}{e^{(p/k_B T_\nu(z))} + 1} \right)$$

$$T_\nu(z) = 1.95(1+z) K$$



$10^{12} h^{-1} M_\odot$

$\sim 100 \text{ km/s}$

$10^{13} h^{-1} M_\odot$

$\sim 200 \text{ km/s}$

$10^{14} h^{-1} M_\odot$

$\sim 450 \text{ km/s}$

$10^{15} h^{-1} M_\odot$

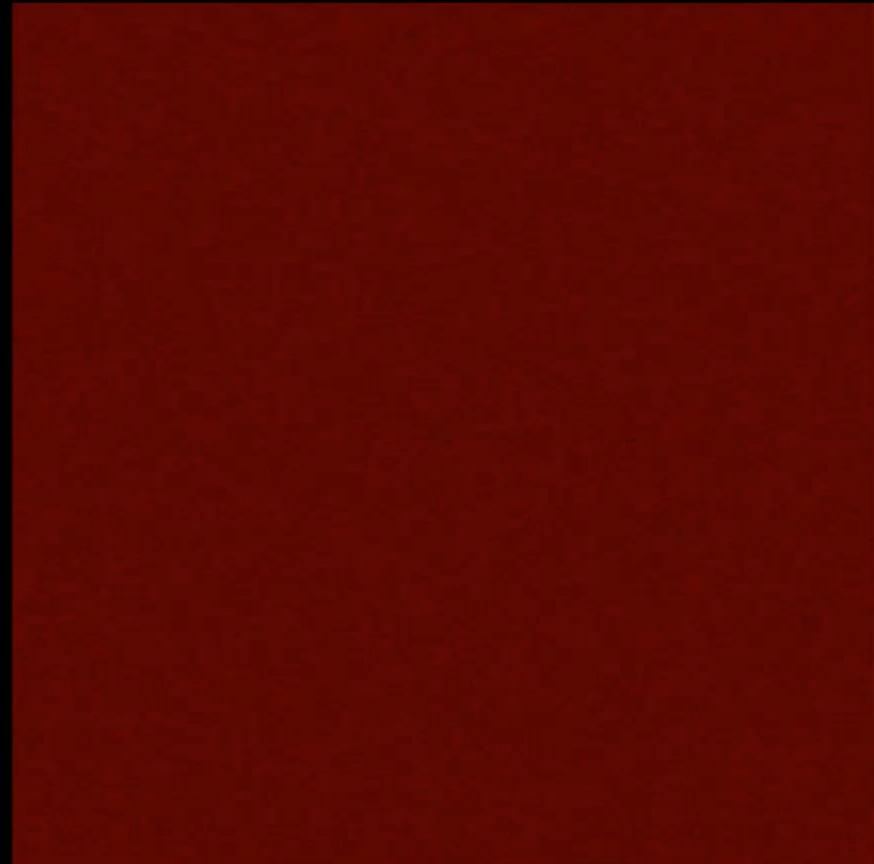
$\sim 950 \text{ km/s}$

Neutrino clustering

Dark Matter

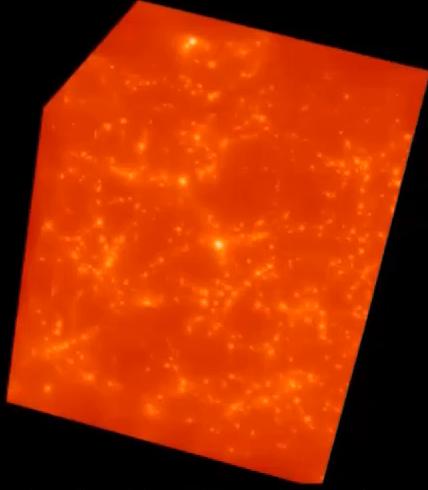


Neutrino

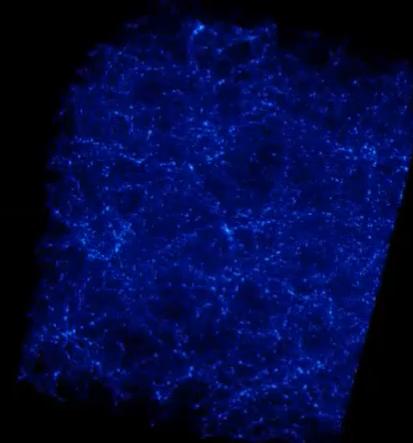


$a=0.02$

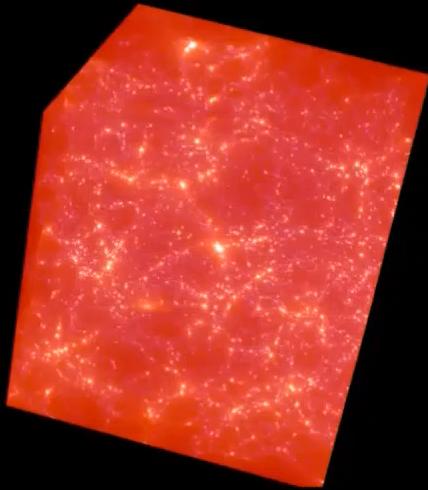
Neutrino



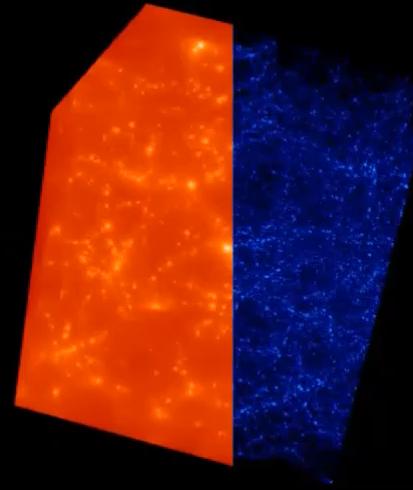
Dark Matter



Blending Neutrino and Dark Matter



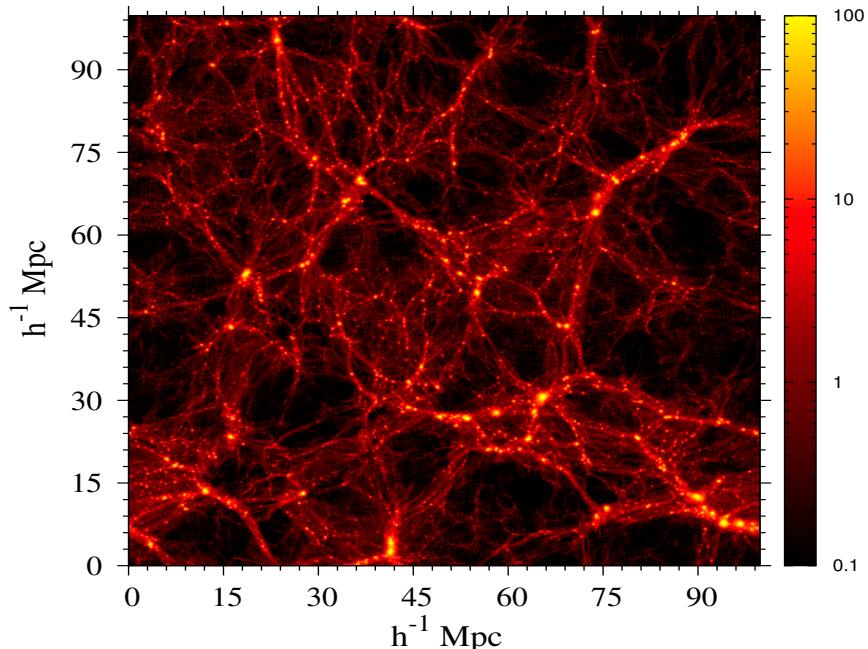
Cropping Neutrino and Dark Matter



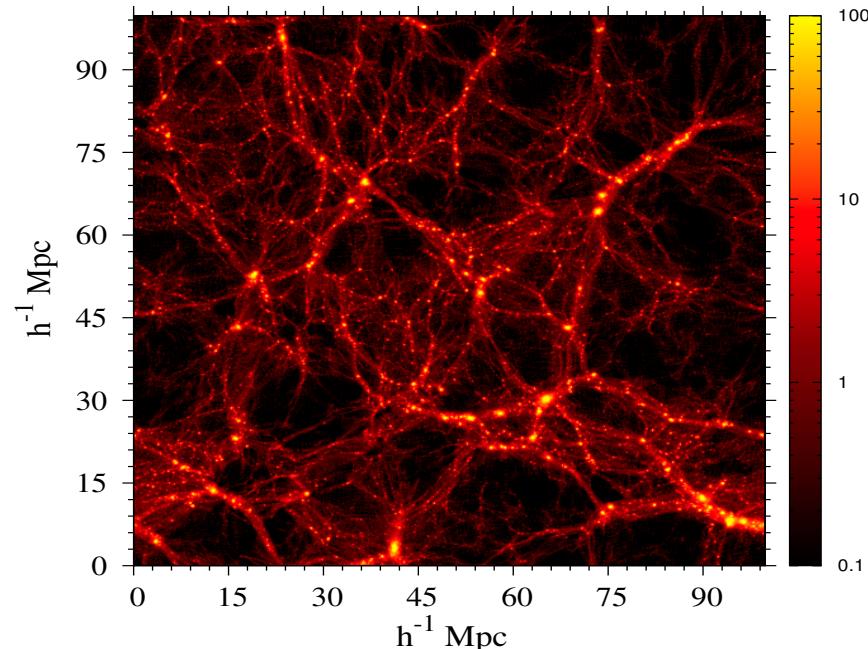
https://www.youtube.com/watch?v=_-Q3dO1lq4g&list=PLBc4nJxplePXY0wdpHfKZi7jHJUjRTaAz

Λ CDM

$$\sum_i m_{\nu_i} = 0.60 \text{ eV}$$



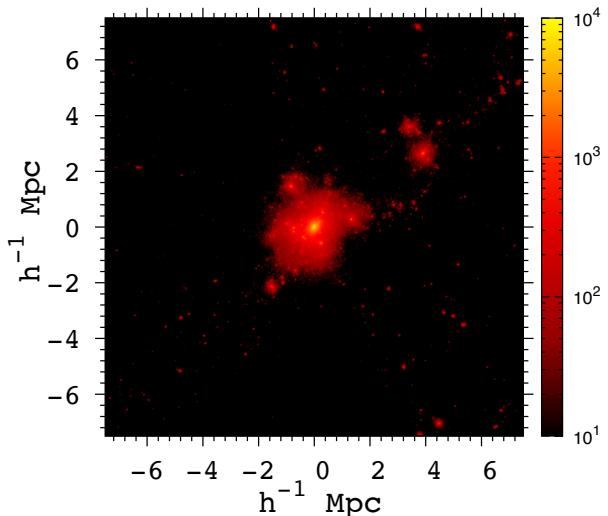
$$\sum_i m_{\nu_i} = 0.30 \text{ eV}$$



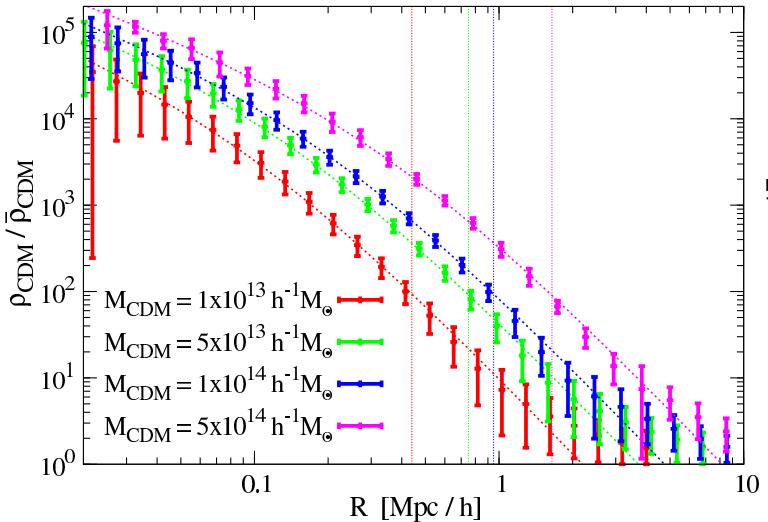
ν

Clustering of relic neutrinos

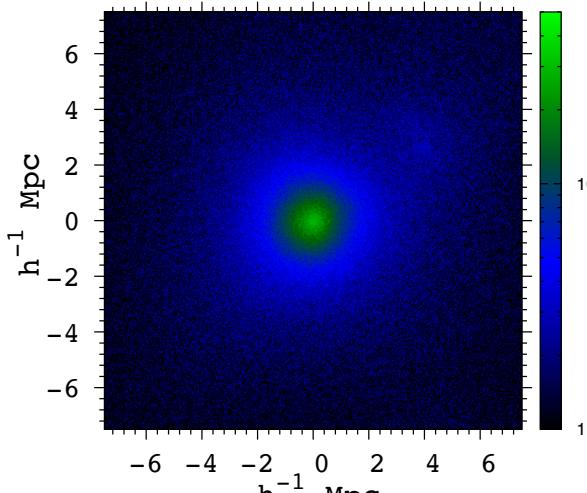
$$M_{CDM} = 4 \times 10^{14} M_\bullet / h$$



$$\rho_{CDM}(r) = \frac{\rho_s}{(r/r_s)(1+r/r_s)^2}$$



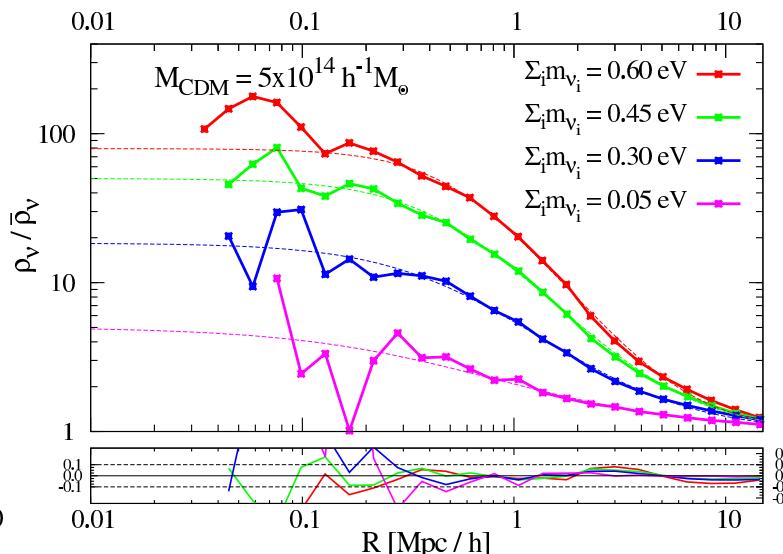
FVN, Bird, Peña-Garay, Viel, 2013
FVN, Miralda-Escude, Peña-Garay, Quilis, 2011



$$C = \frac{r_\nu}{r_s}$$

$$\delta_\nu(r) = \frac{\rho_c}{(1+r/r_c)^\alpha}$$

$$F_h = \begin{cases} 9.5 \times 10^{-4} \rightarrow 0.3 \text{ eV} \\ 2.6 \times 10^{-3} \rightarrow 0.6 \text{ eV} \end{cases}$$



PTOLEMY

Princeton Tritium Observatory
for Light, Early-Universe,
Massive-Neutrino Yield

Halo mass function

Castorina, Sefussati, Sheth, FVN, Viel 2013

$$\frac{dn(M,z)}{dM} = \nu f(\nu) \frac{\rho_m}{M^2} \frac{d \ln \nu}{d \ln M}$$

Universal

$\nu \equiv \frac{\delta_c}{\sigma(M,z)} \quad \delta_c = 1.686$ $\sigma^2(M,z) = \frac{1}{2\pi^2} \int_0^\infty k^2 P_m(k) W^2(k,R) dk$ $M = \frac{4\pi}{3} \rho_m R^3$

What about massive neutrino cosmologies?

- No ~~prescription~~

Brandbyge et al. 2010

$$\rho_m \rightarrow \cancel{\rho_m} \rightarrow \cancel{P_m(k)} \rightarrow P_m(k)$$

- Matter prescription

Brandbyge et al. 2010
Marulli et al. 2011
Villaescusa-Navarro et al. 2013

$$\rho_m \rightarrow \rho_{cdm} \quad P_m(k) \rightarrow P_m(k)$$

- Cold dark matter prescription

Ichiki & Takada 2011
Castorina et al. 2013
Costanzi et al. 2013

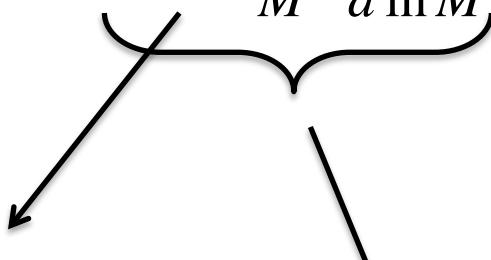
$$\rho_m \rightarrow \rho_{cdm} \quad P_m(k) \rightarrow P_{cdm}(k)$$

Halo mass function

Castorina, Sefussati, Sheth, FVN, Viel 2013

FoF halos : $b=0.2$

$$\frac{dn(M, z)}{dM} = \nu f(\nu) \frac{\rho_m}{M^2} \frac{d \ln \nu}{d \ln M}$$



Crocce et al. 2010

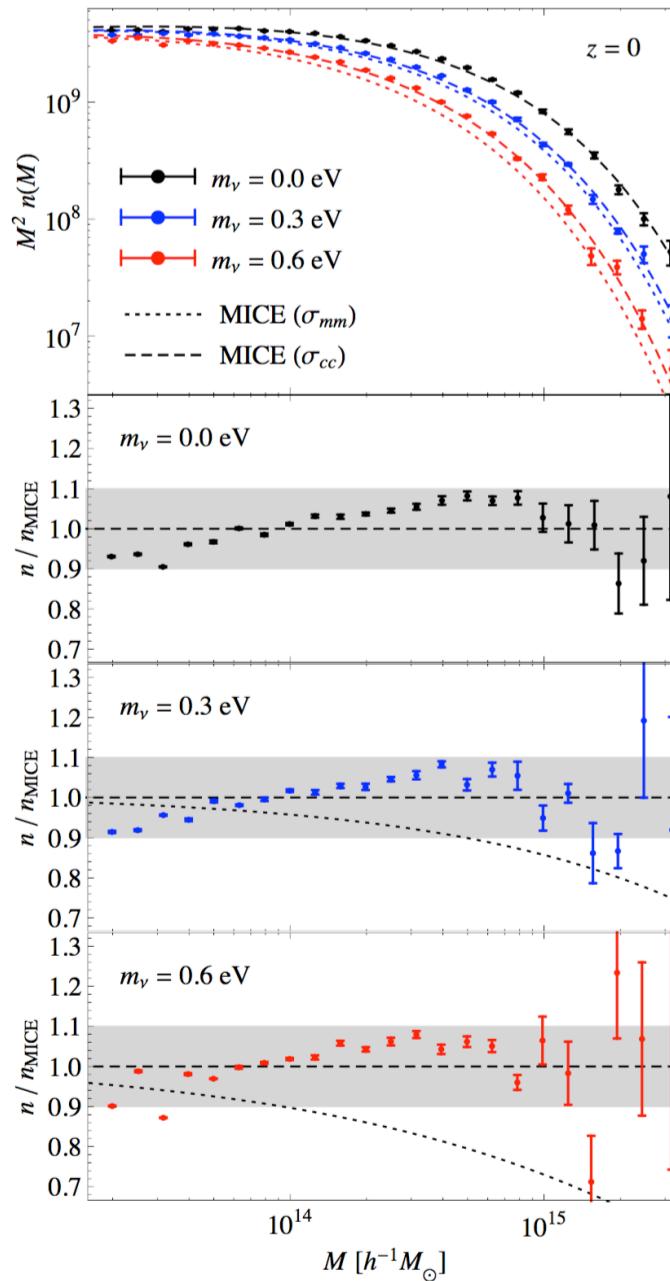
Matter prescription

$$\rho_m \rightarrow \rho_{cdm} \quad P_m(k) \rightarrow P_m(k)$$



Cold dark matter prescription

$$\rho_m \rightarrow \rho_{cdm} \quad P_m(k) \rightarrow P_{cdm}(k)$$



Halo mass function

Castorina, Sefussati, Sheth, FVN, Viel 2013

FoF halos : $b=0.2$

$$\frac{dn(M, z)}{dM} = \nu f(\nu) \frac{\rho_m}{M^2} \frac{d \ln \nu}{d \ln M}$$

Universal?

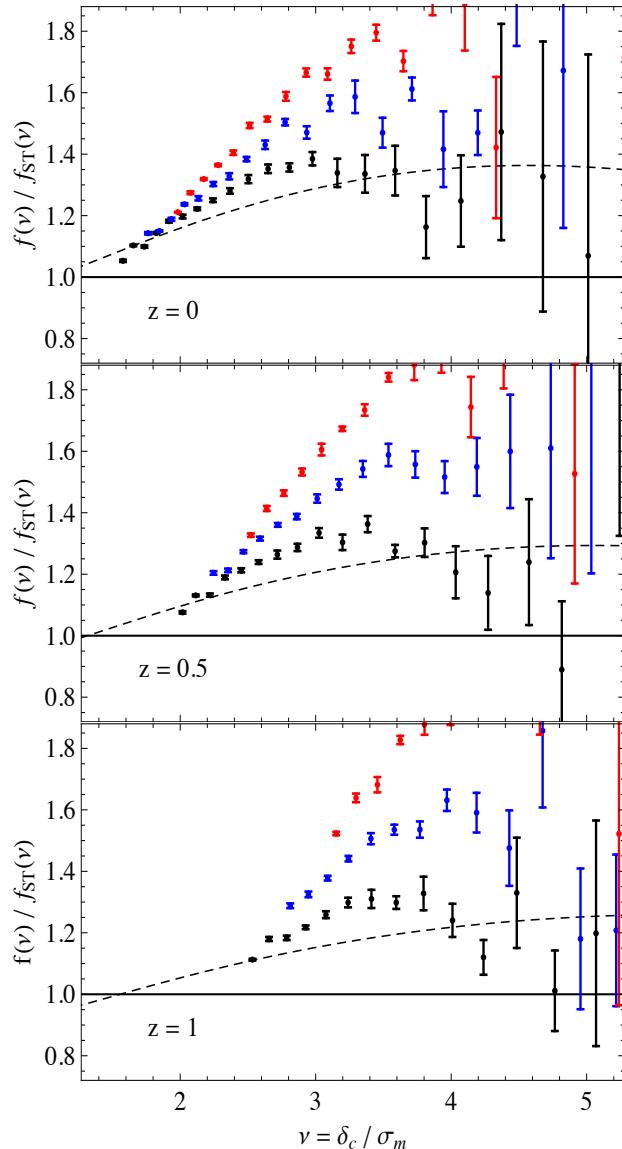
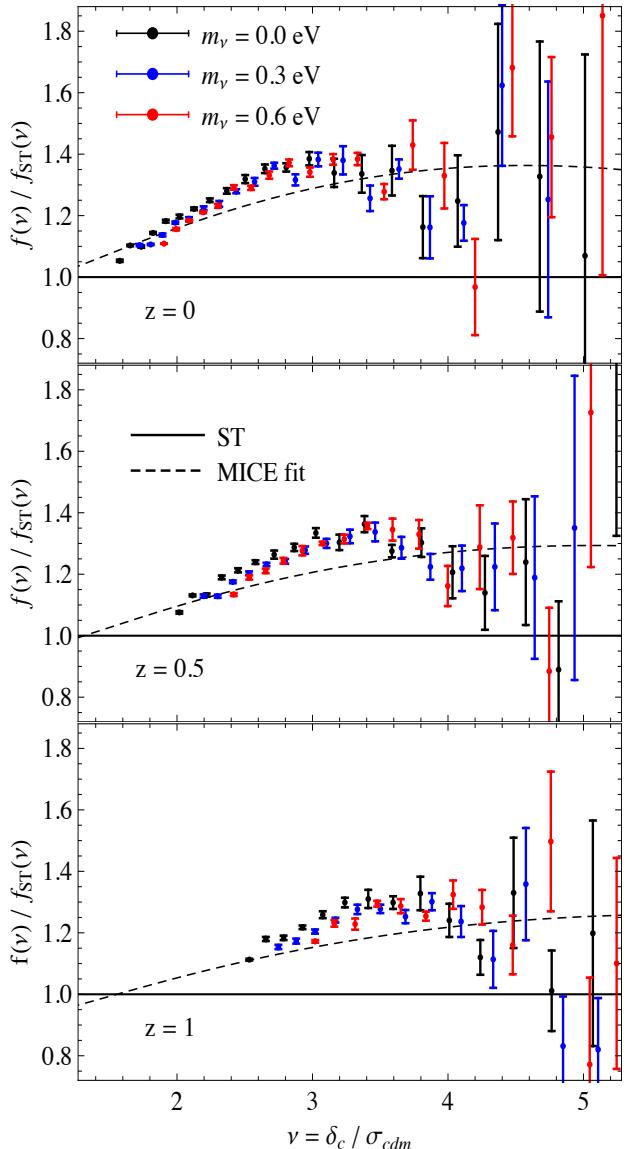
$$f(\nu) = \frac{M^2}{\rho} \frac{1}{\nu} \frac{d \ln M}{d \ln \nu} \frac{dn(M, z)}{dM}$$

Matter ~~prescription~~

$$\rho_m \rightarrow \rho_{cdm} \quad P_m(k) \rightarrow P_{cdm}(k)$$

Cold dark matter prescription

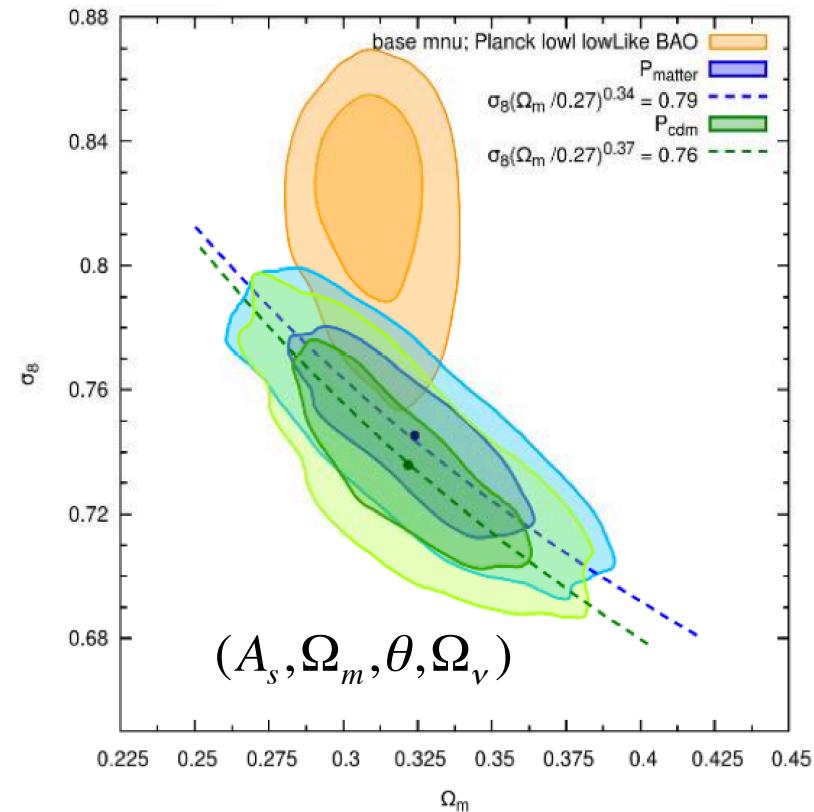
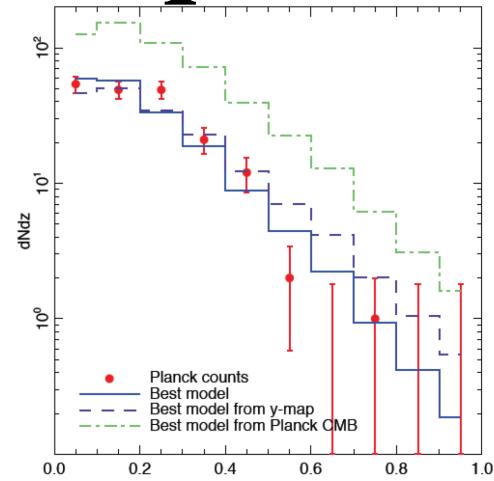
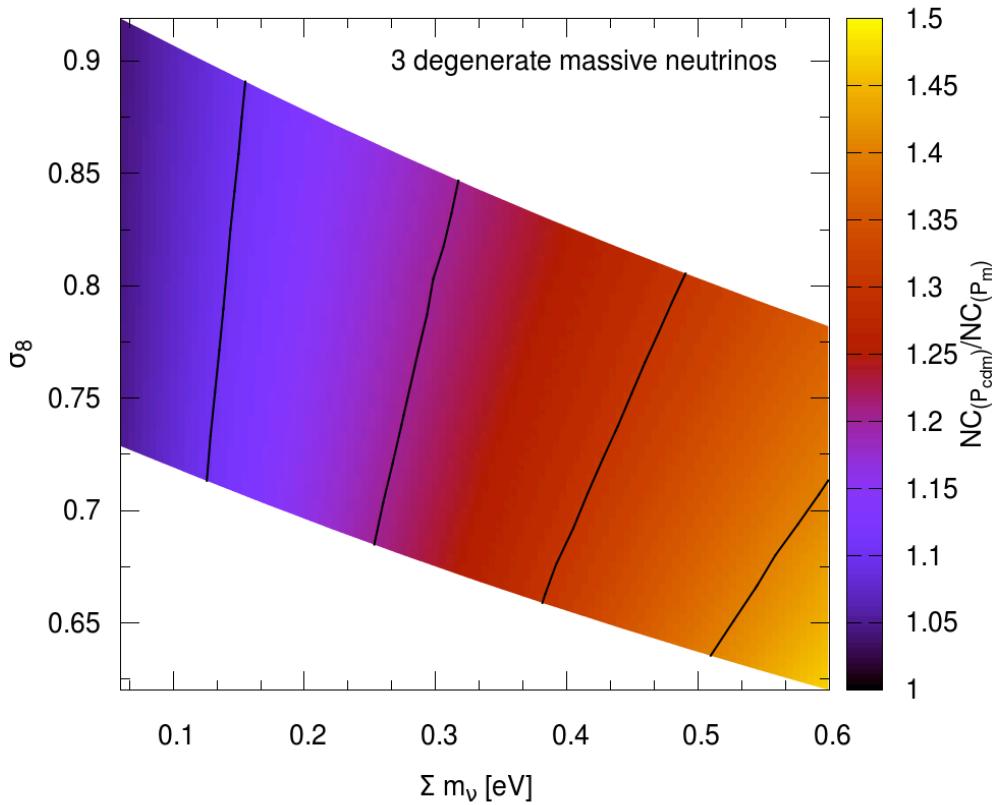
$$\rho_m \rightarrow \rho_{cdm} \quad P_m(k) \rightarrow P_{cdm}(k)$$



Halo mass function: consequences

Costanzi, FVN, Viel, Xia, Borgani, Castorina, Sefusatti 2013

$$N_i = \int_{z_i}^{z_{i+1}} dz \int_{\Delta\Omega} d\Omega \frac{dV}{dz d\Omega} \int_0^\infty dM X(M, z, \Omega) n(M, z),$$



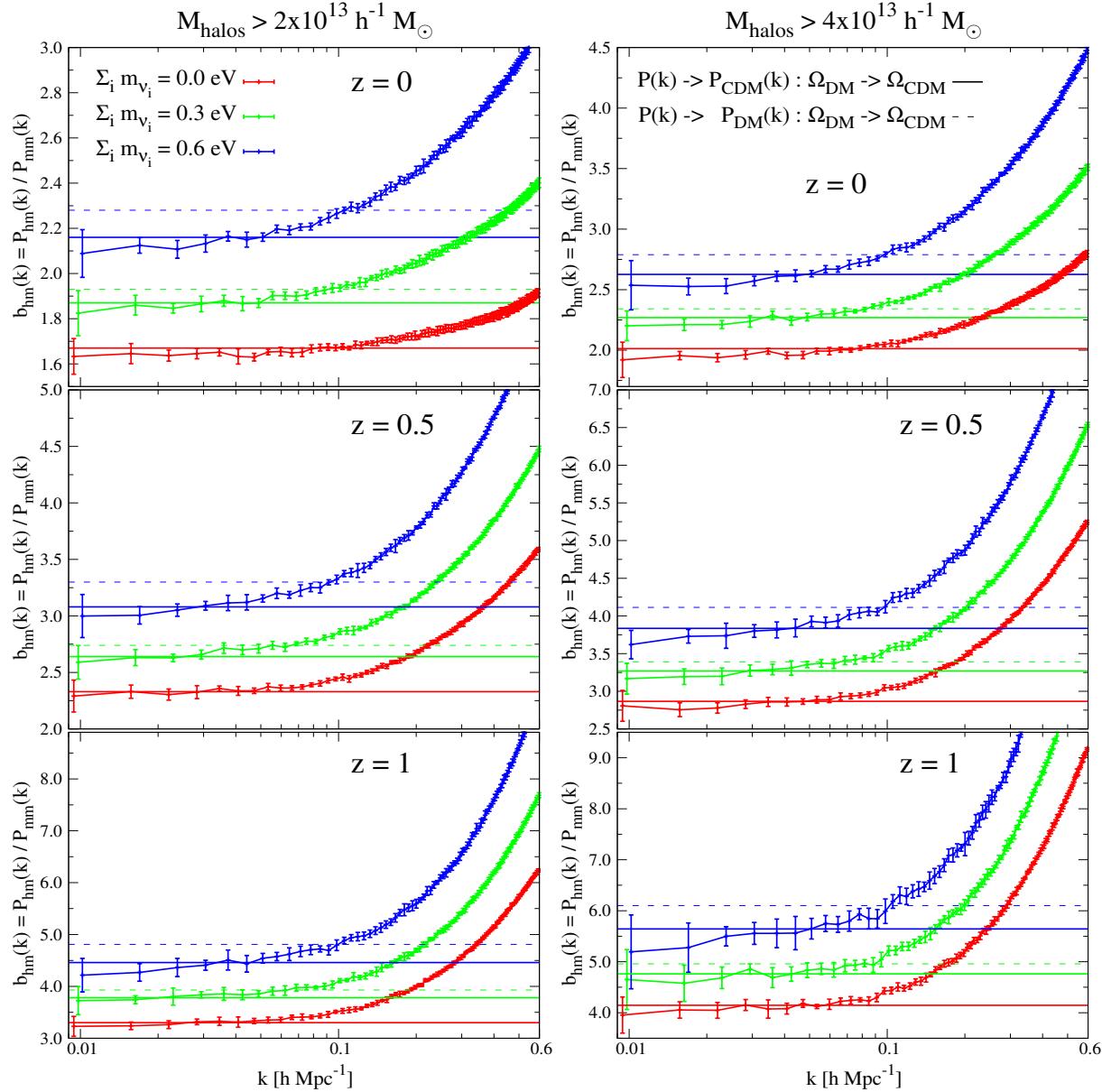
Clustering of dark matter halos

FVN, Marulli, Viel, Branchini, Castorina, Sefusatti, Saito 2013

Castorina, Sefusatti, Sheth, FVN, Viel 2013

$$b_{hm}(k) = \frac{P_{hm}(k)}{P_{mm}(k)}$$

- 0.0 eV → 8 realizations
- 0.3 eV → 8 realizations
- 0.6 eV → 8 realizations



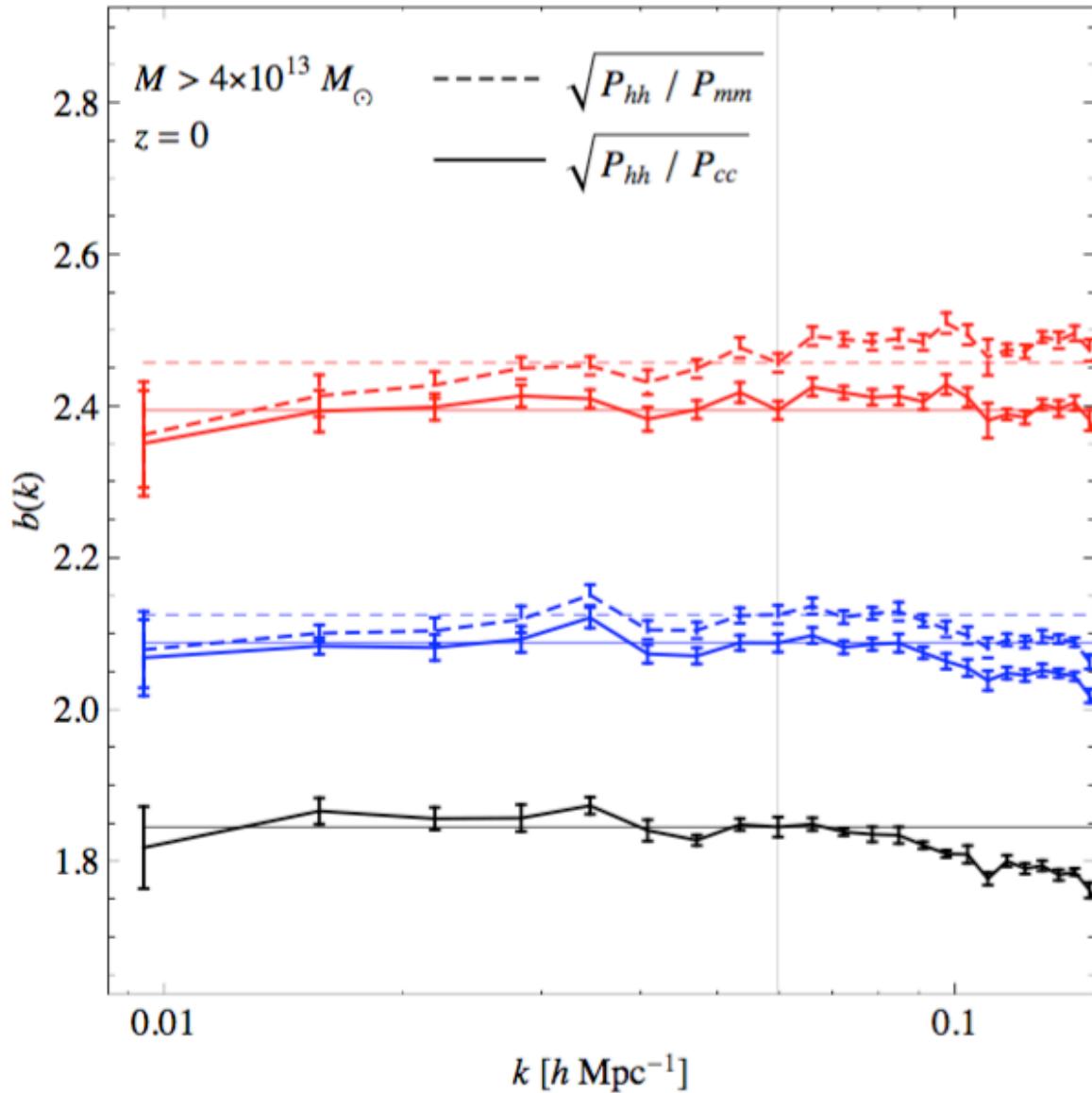
Clustering of dark matter halos

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FVN, Marulli, Viel, Branchini, Castorina, Sefusatti, Saito 2013

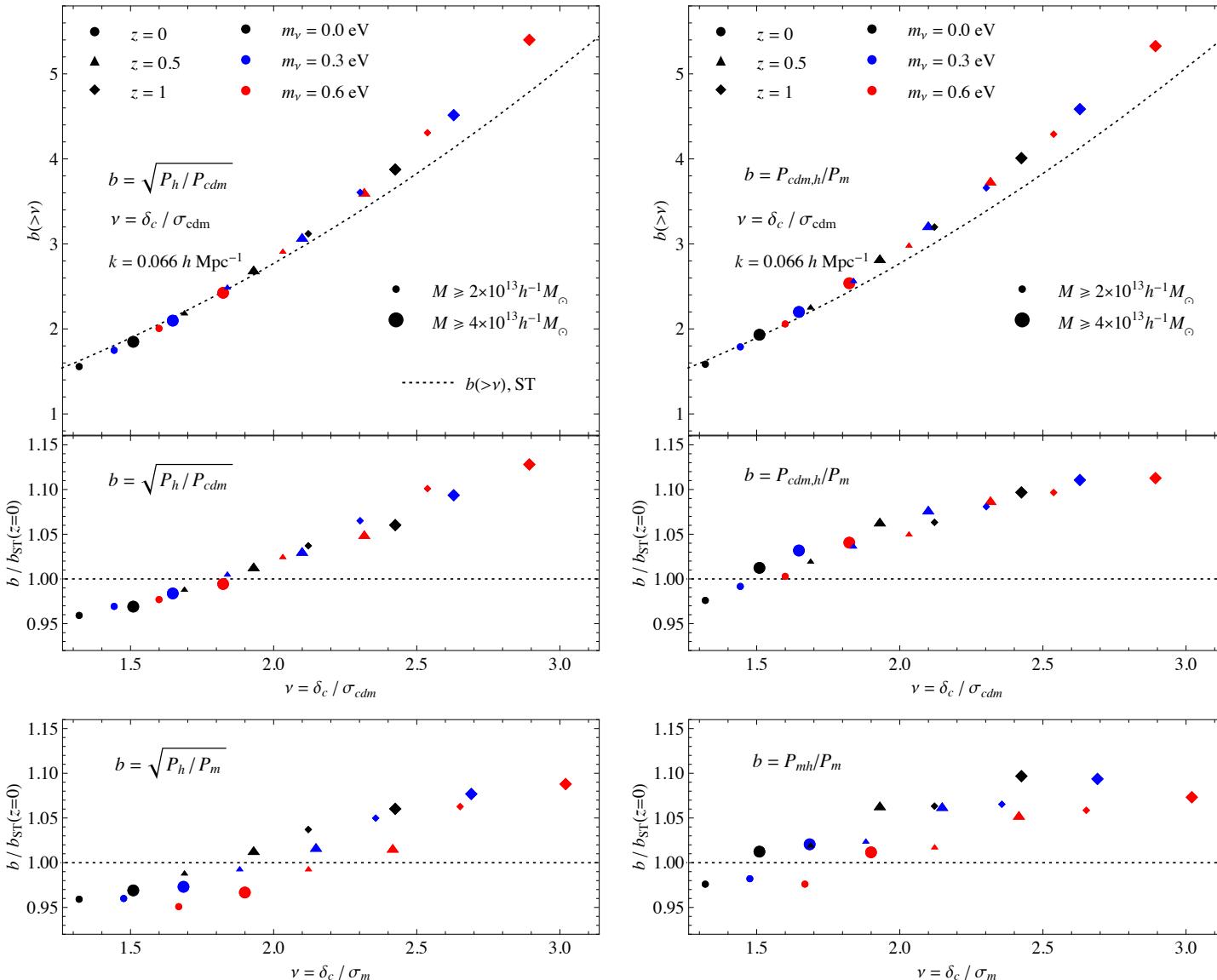
$$b_{hm}(k) = \frac{P_{hm}(k)}{P_{mm}(k)}$$

$$b_{hc}(k) = \frac{P_{hc}(k)}{P_{cc}(k)}$$



Clustering of dark matter halos

Castorina, Sefusatti, Sheth, FVN, Viel 2013



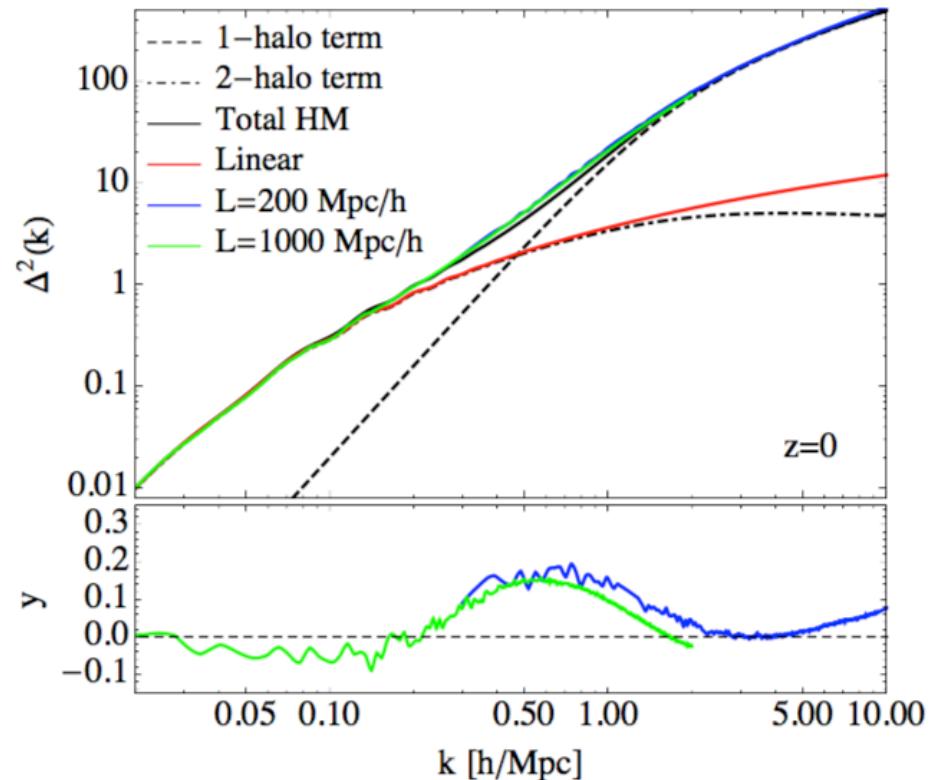
Halo model

Seljak 2000; Ma & Fry 2000
Peacock & Smith 2000; Cooray & Sheth 2002

Assumptions

- All matter in halos
- Linear matter power spectrum
- Halo mass function
- Halo bias
- Density profile

Ingredients



$$P(k) = P_{1h}(k) + P_{2h}(k)$$

$$P_{1h}(k) = \int_0^\infty f(\nu) \frac{M}{\rho} |u(k|\nu)|^2 d\nu$$

$$P_{2h}(k) = \left[\int_0^\infty b(\nu) f(\nu) u(k|\nu) d\nu \right]^2 P_L(k)$$

Halo model

Massara, FVN, Viel 2014

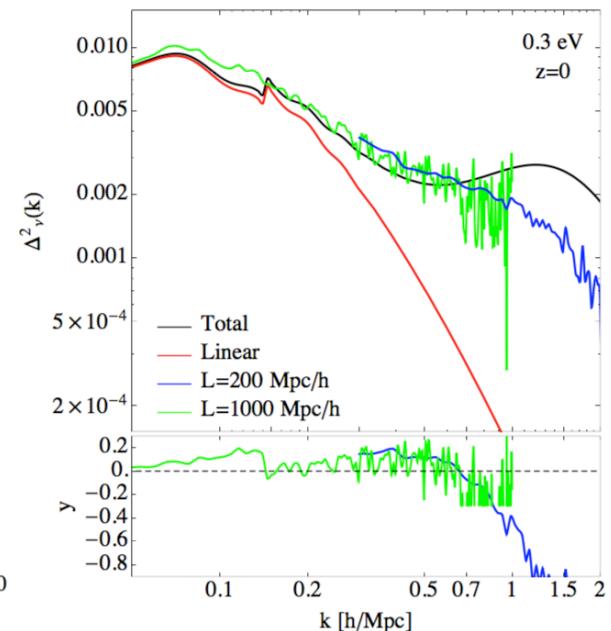
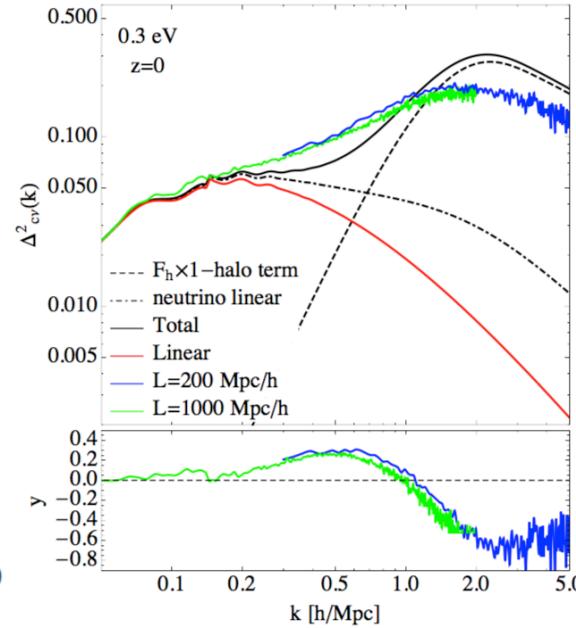
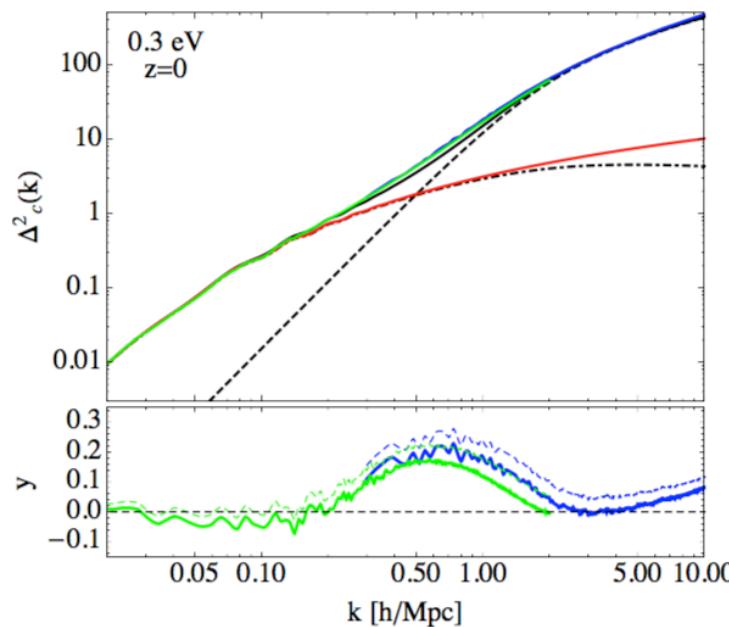
$$\rho_m = \rho_{cb} + \rho_\nu$$

$$P_m(k) = \left(\frac{\Omega_{cb}}{\Omega_m} \right)^2 P_{cb}(k) + 2 \left(\frac{\Omega_{cb} \Omega_\nu}{\Omega_m} \right) P_{cb\nu}(k) + \left(\frac{\Omega_\nu}{\Omega_m} \right)^2 P_\nu(k)$$

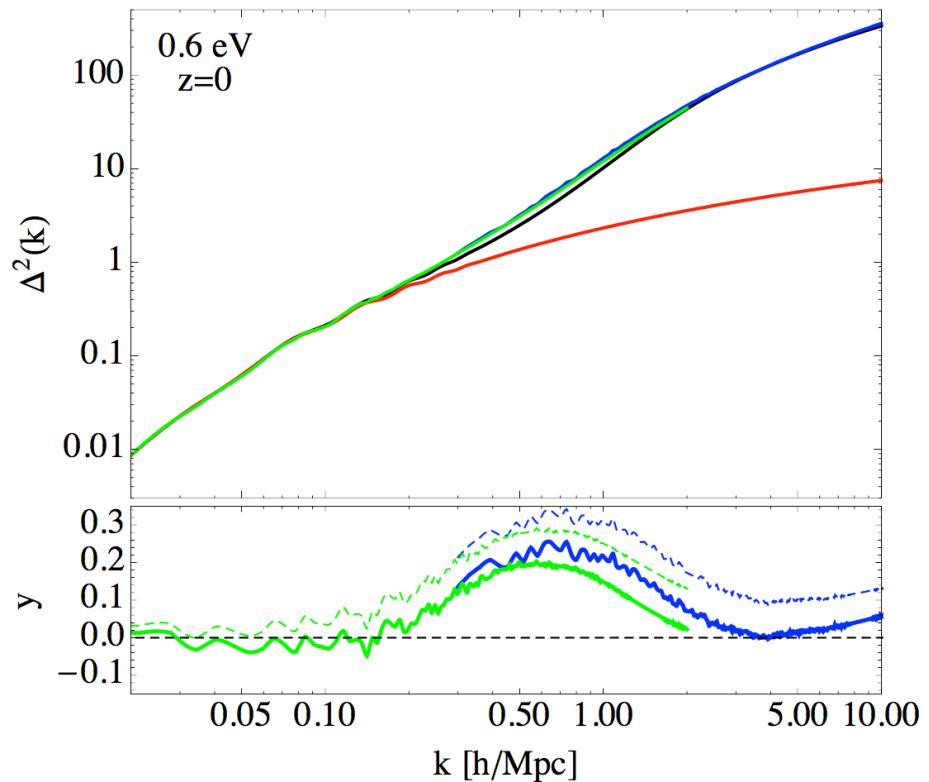
$$\delta_\nu = F_h \delta_\nu^h + (1 - F_h) \delta_\nu^h$$

$$\delta_\nu(r) = \frac{\rho_c}{(1 + r/r_c)^\alpha}$$

- All matter in halos
- HMF and halo bias as a function of matter field
- All CDM in halos
- HMF and halo bias as a function of CDM field
- Concentration in halos
- Only a fraction of ν in halos

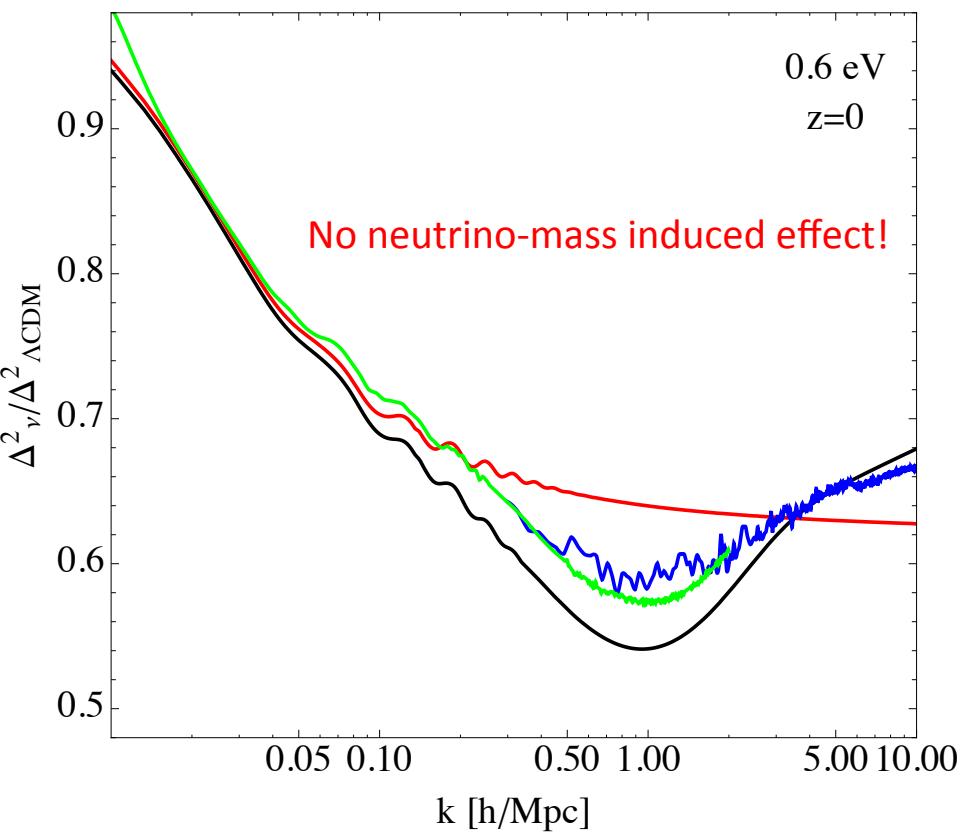


Halo model



Massara, FVN, Viel 2014

$$\frac{P_{\nu\Lambda\text{CDM}}(k)}{P_{\Lambda\text{CDM}}(k)} \left\{ \begin{array}{l} M_\nu = 0.15 \text{ eV} \rightarrow 2\% \\ M_\nu = 0.30 \text{ eV} \rightarrow 5\% \\ M_\nu = 0.60 \text{ eV} \rightarrow 10\% \end{array} \right.$$



Neutrino effects on BAO

Peloso, Pietroni, Viel, FVN, 2015

100 N-body simulations

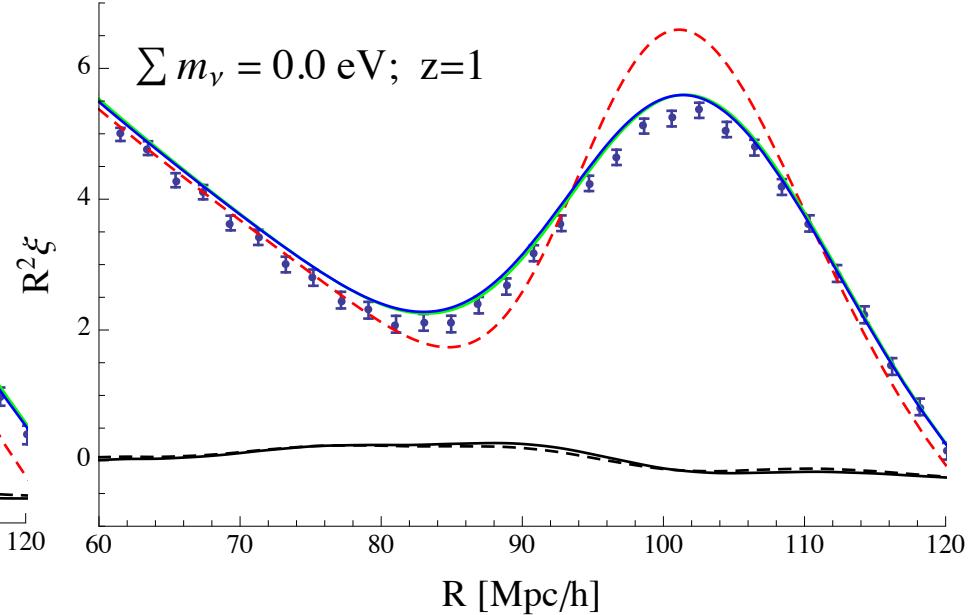
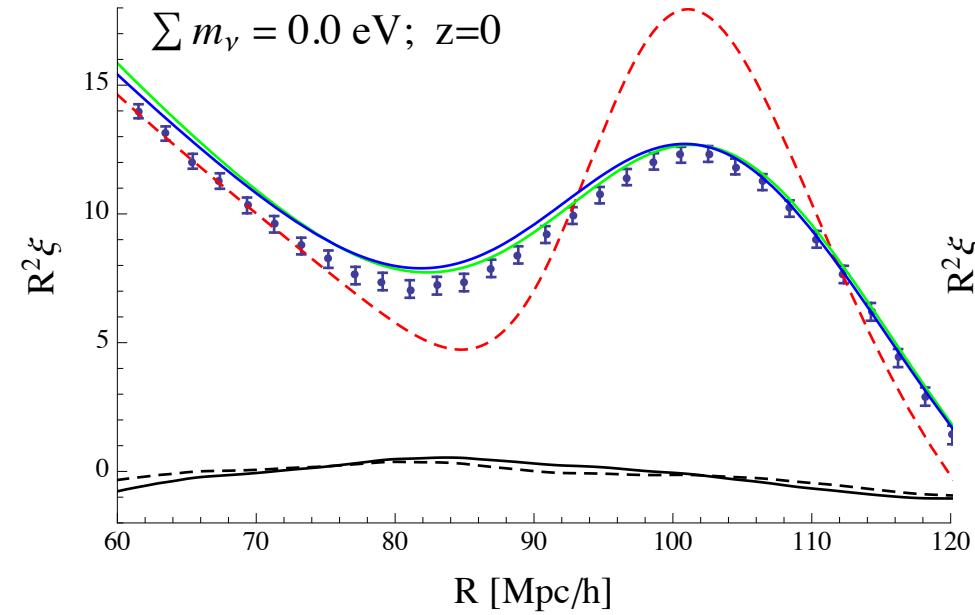
$z_i=99$; 1000 Mpc/h

256^3 CDM + 256^3 Neutrinos

0.0 eV, 0.15 eV, 0.3 eV, 0.6 eV

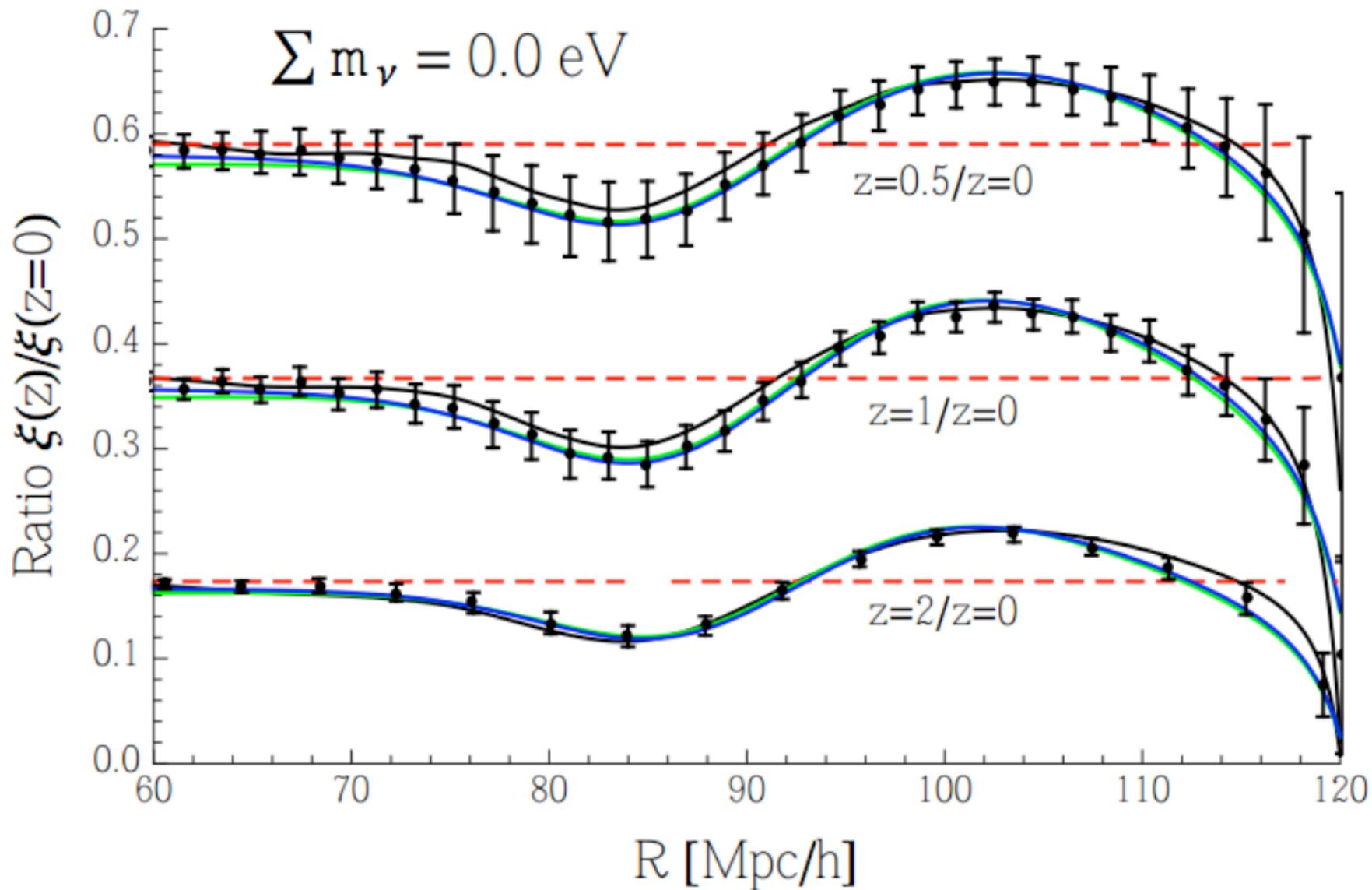
$$P^{(1)}(k, z) = e^{-k^2 \sigma_v^2(z)} P^{lin}(k, z),$$

$$\sigma_v^2(z) = \frac{1}{3} \int \frac{d^3 q}{(2\pi)^3} \frac{P^{lin}(q, z)}{q^2}.$$



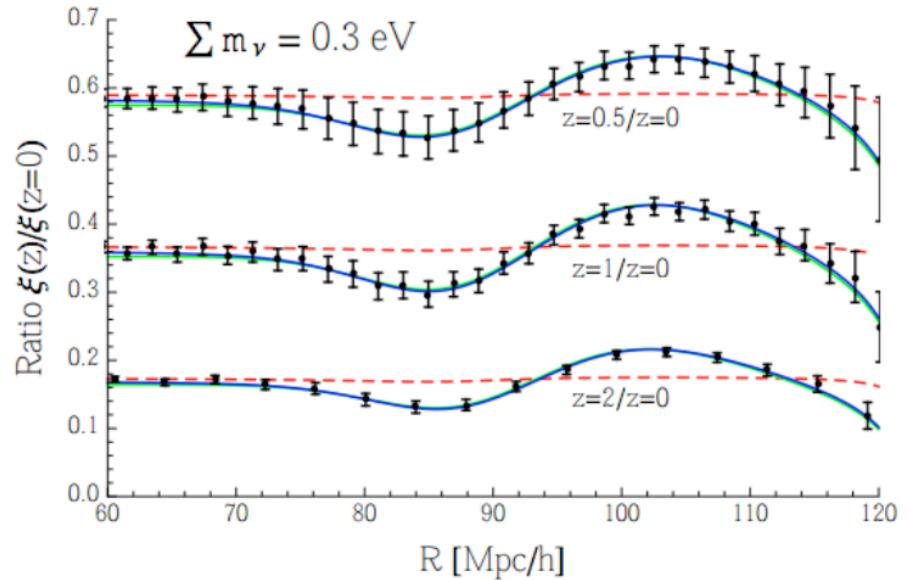
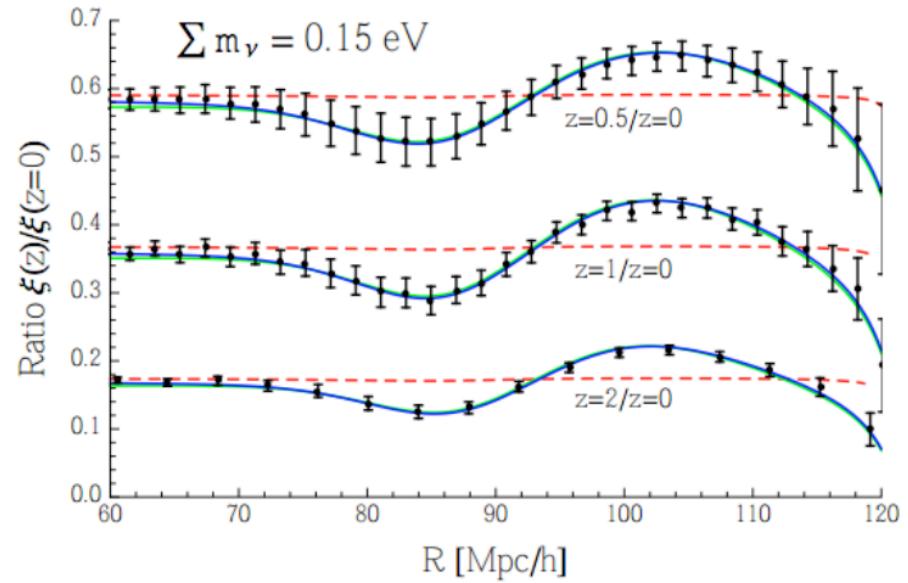
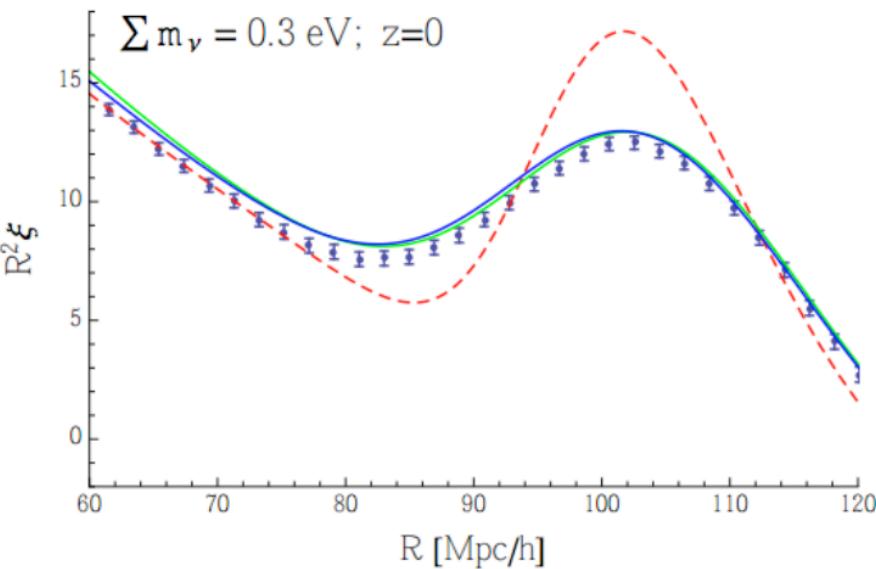
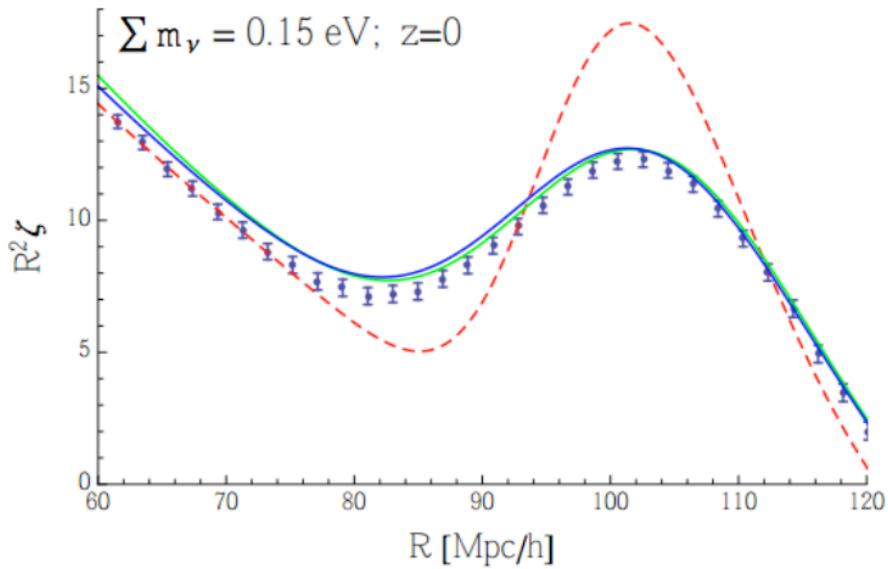
Neutrino effects on BAO

Peloso, Pietroni, Viel, FVN, 2015



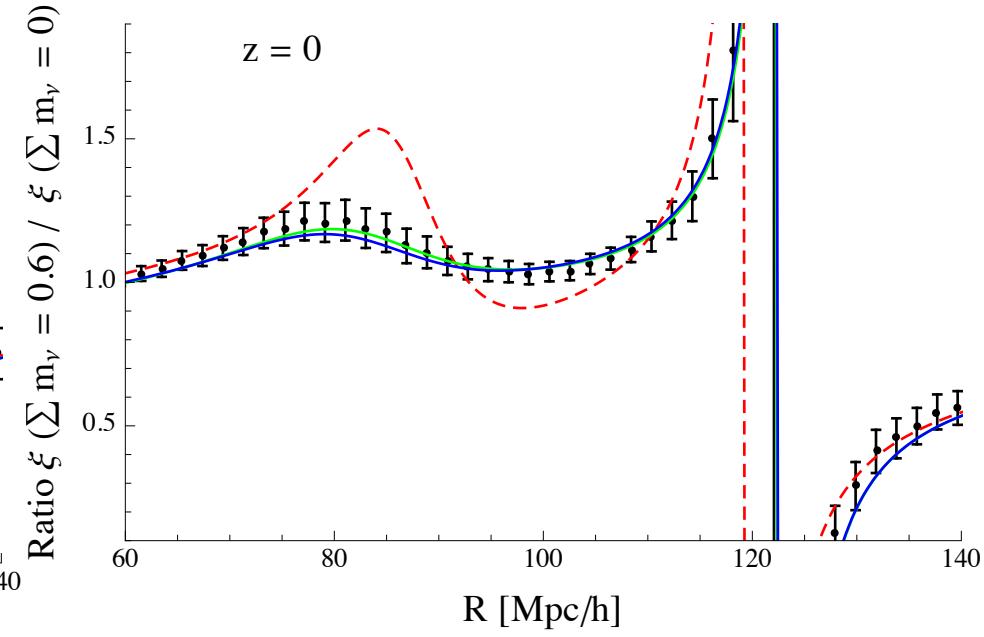
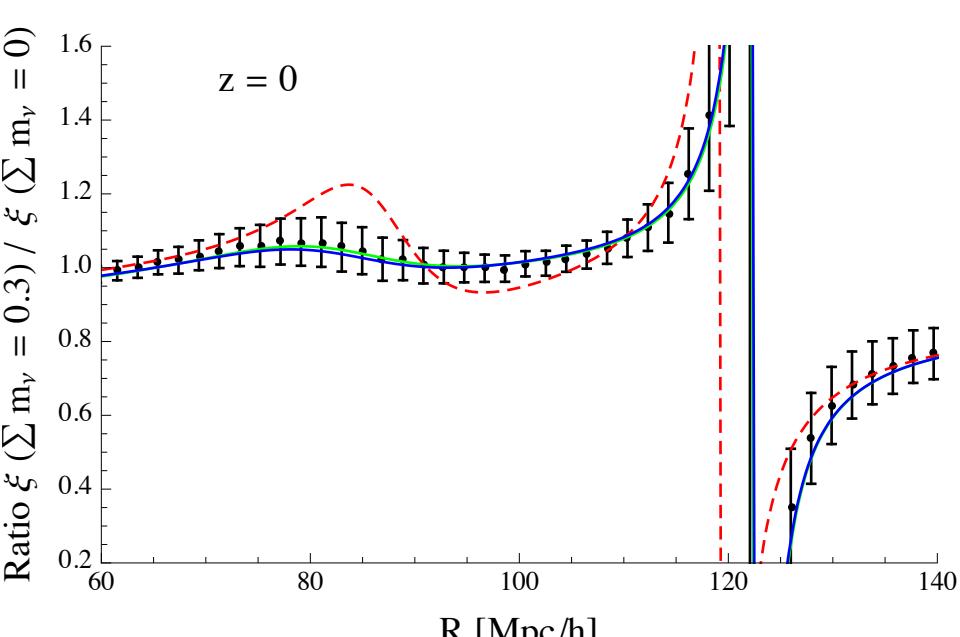
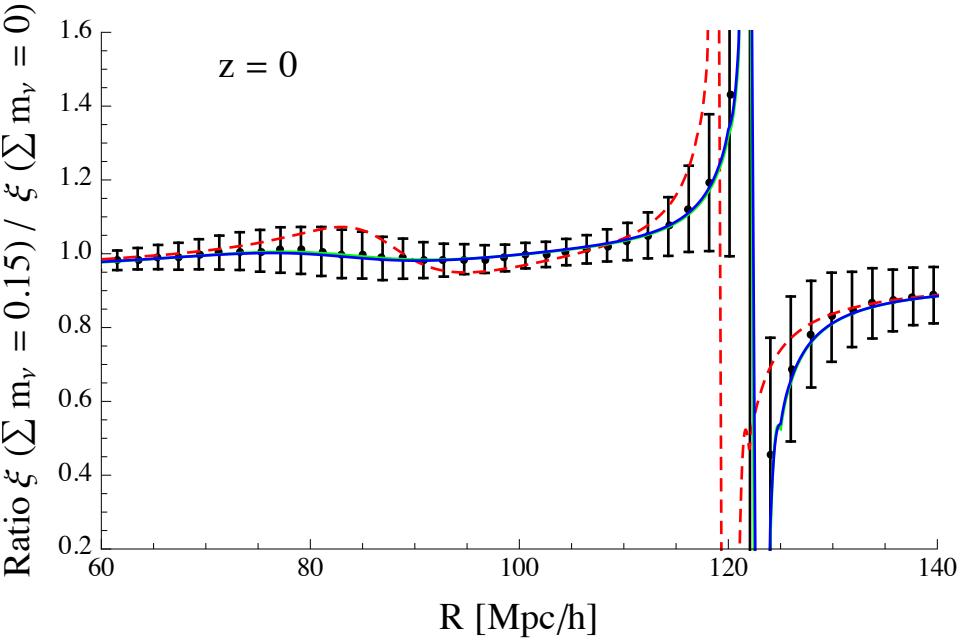
Neutrino effects on BAO

Peloso, Pietroni, Viel, FVN 2015



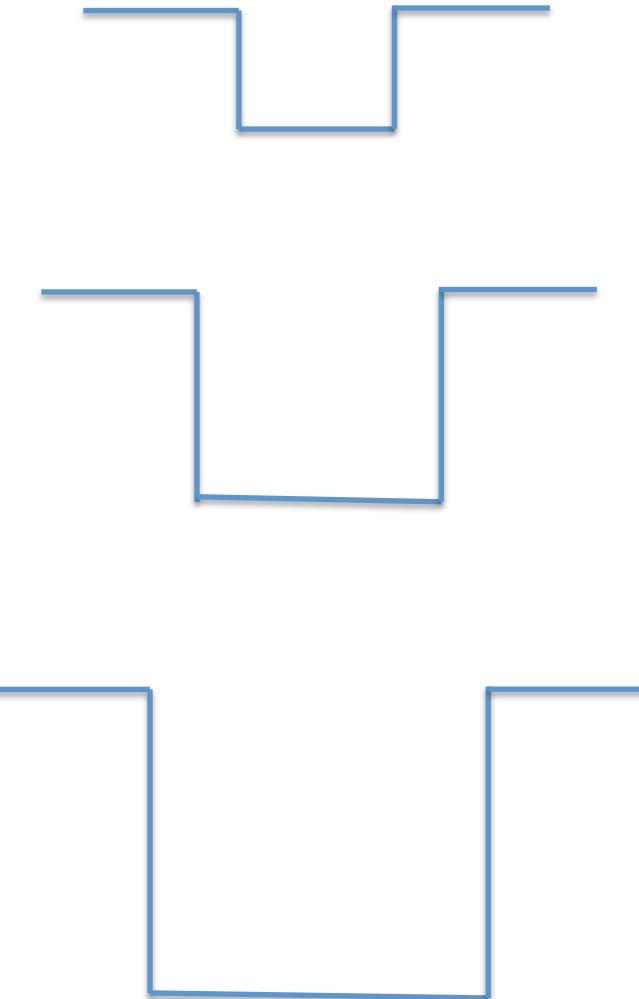
Neutrino effects on BAO

Peloso, Pietroni, Viel, FVN 2015

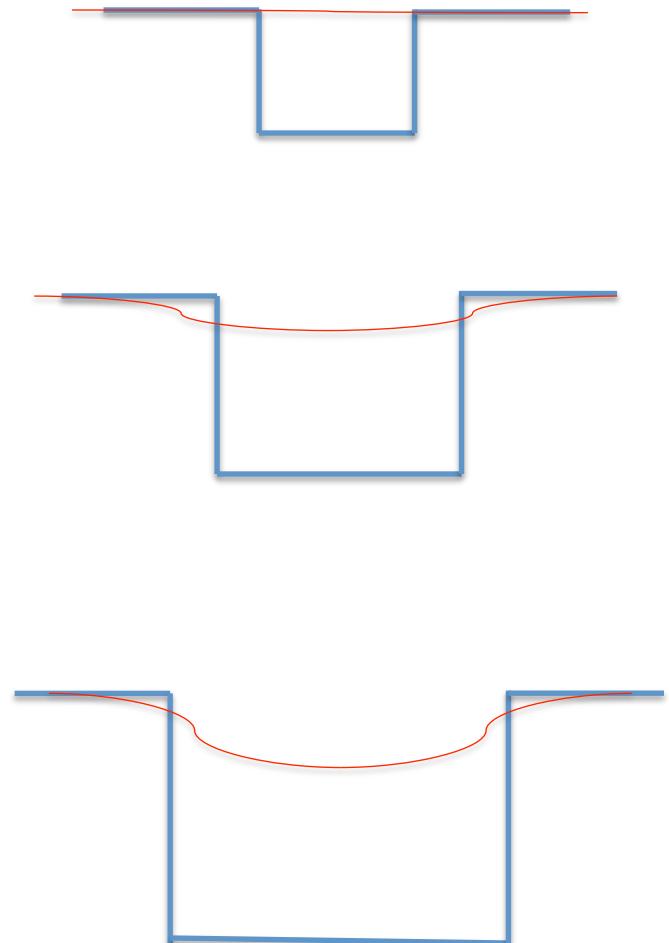


Neutrino effects on voids

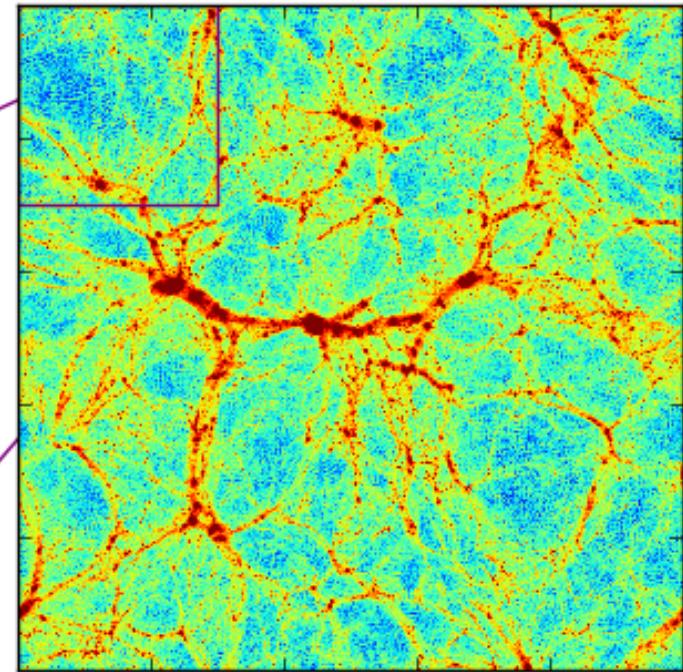
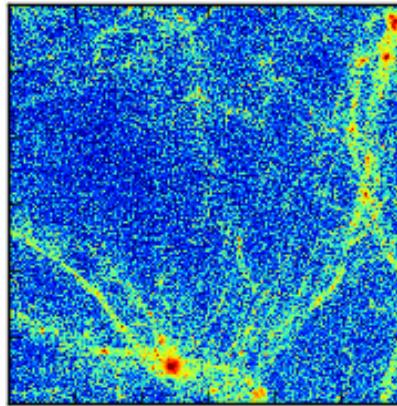
Massless neutrinos



Massive neutrinos

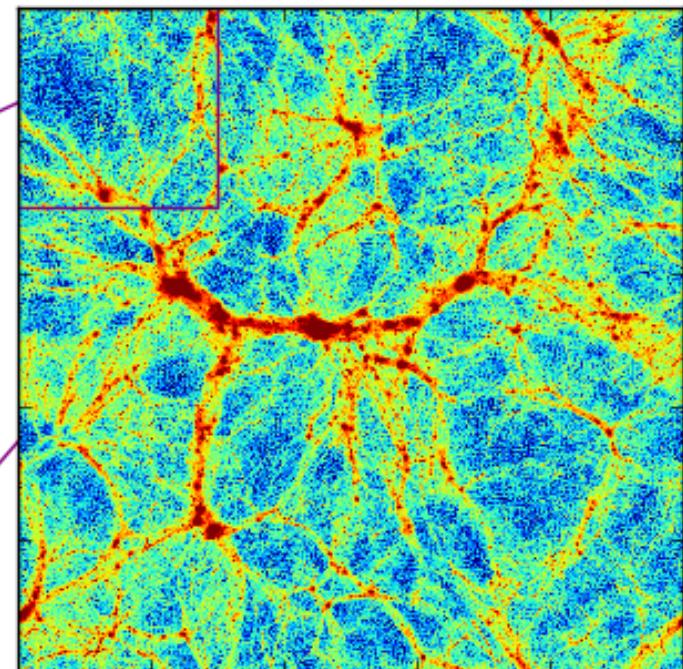
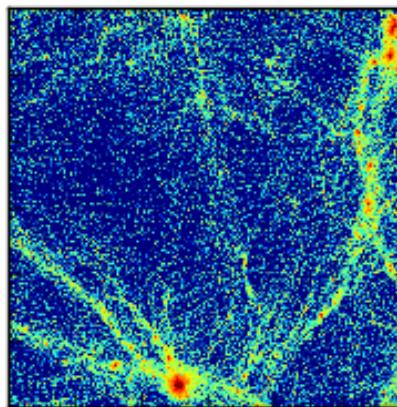


0.6 eV →



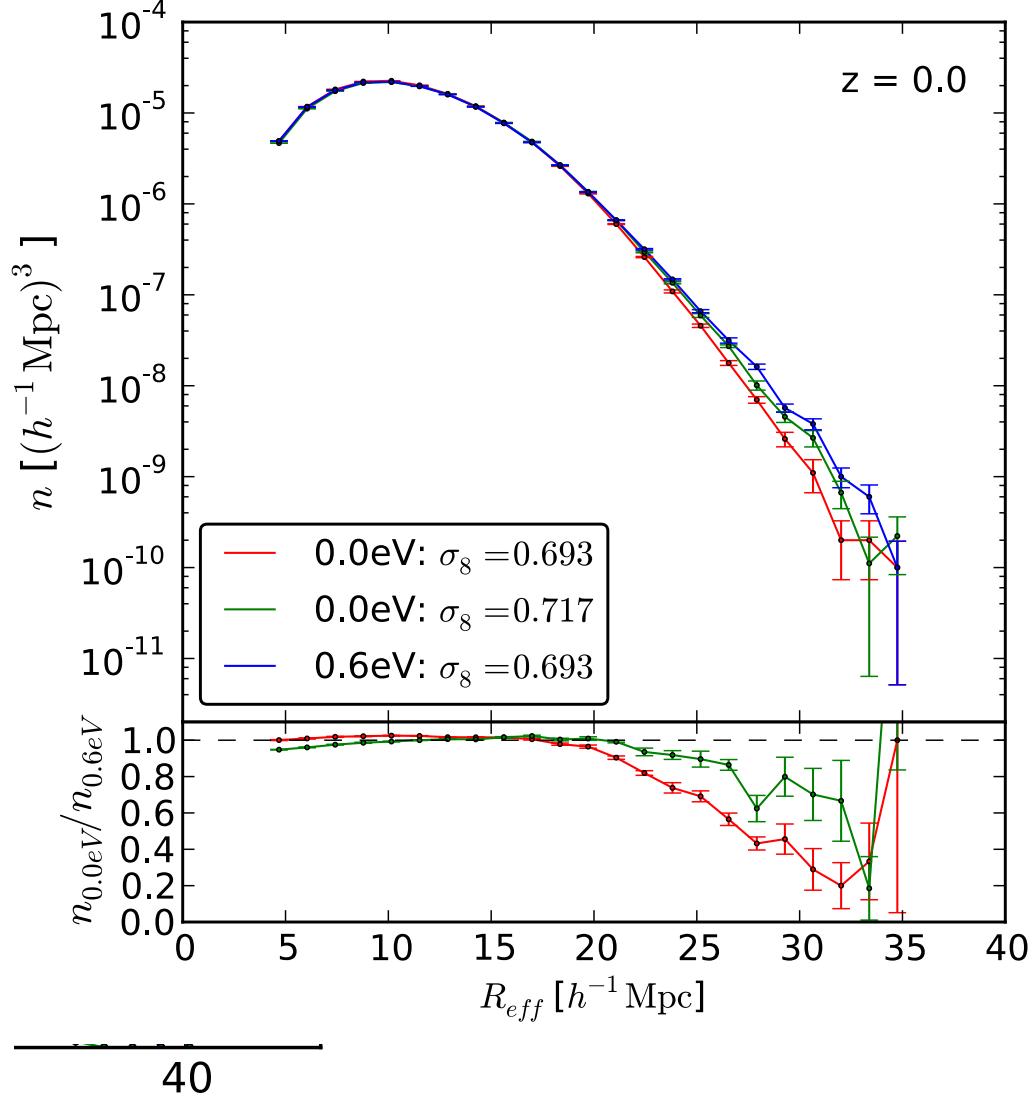
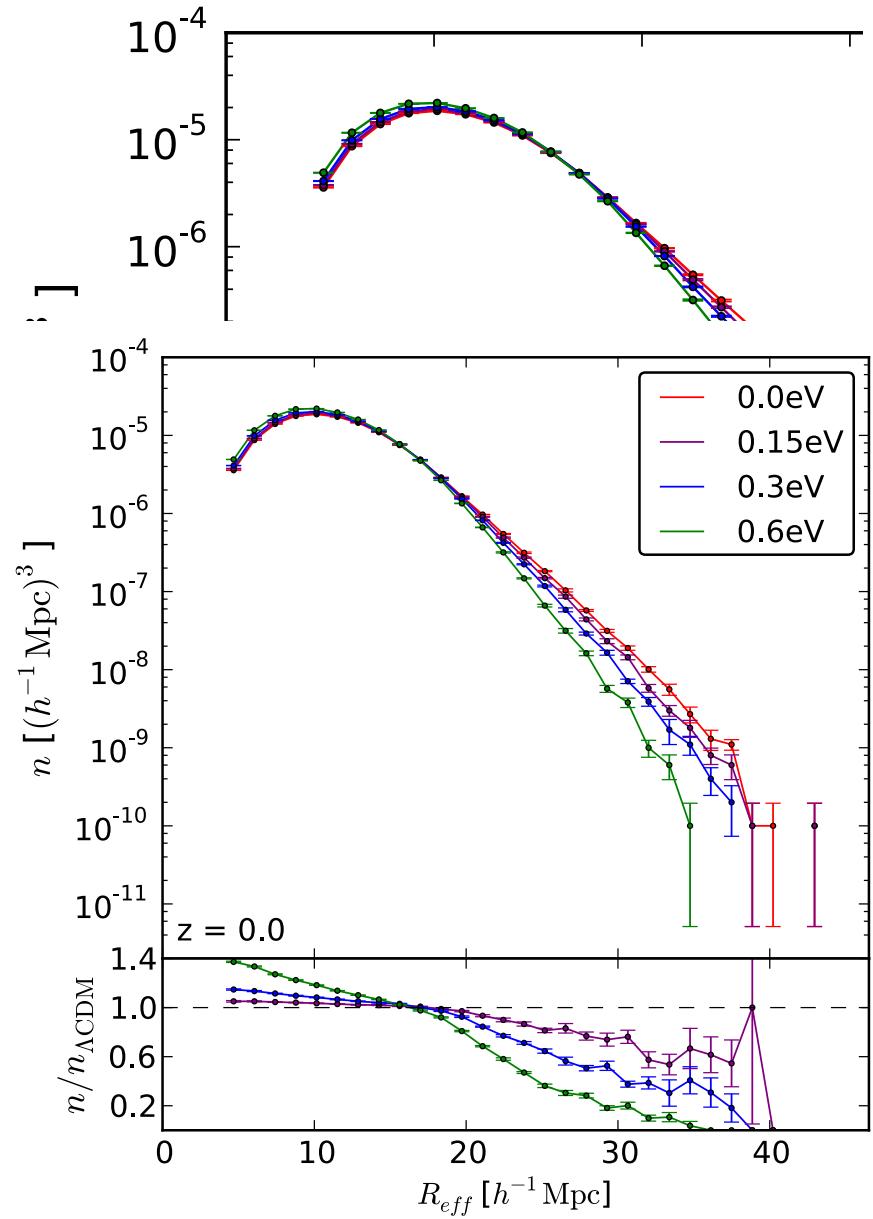
The impact of
massive neutrinos
on cosmic voids

0.0 eV →



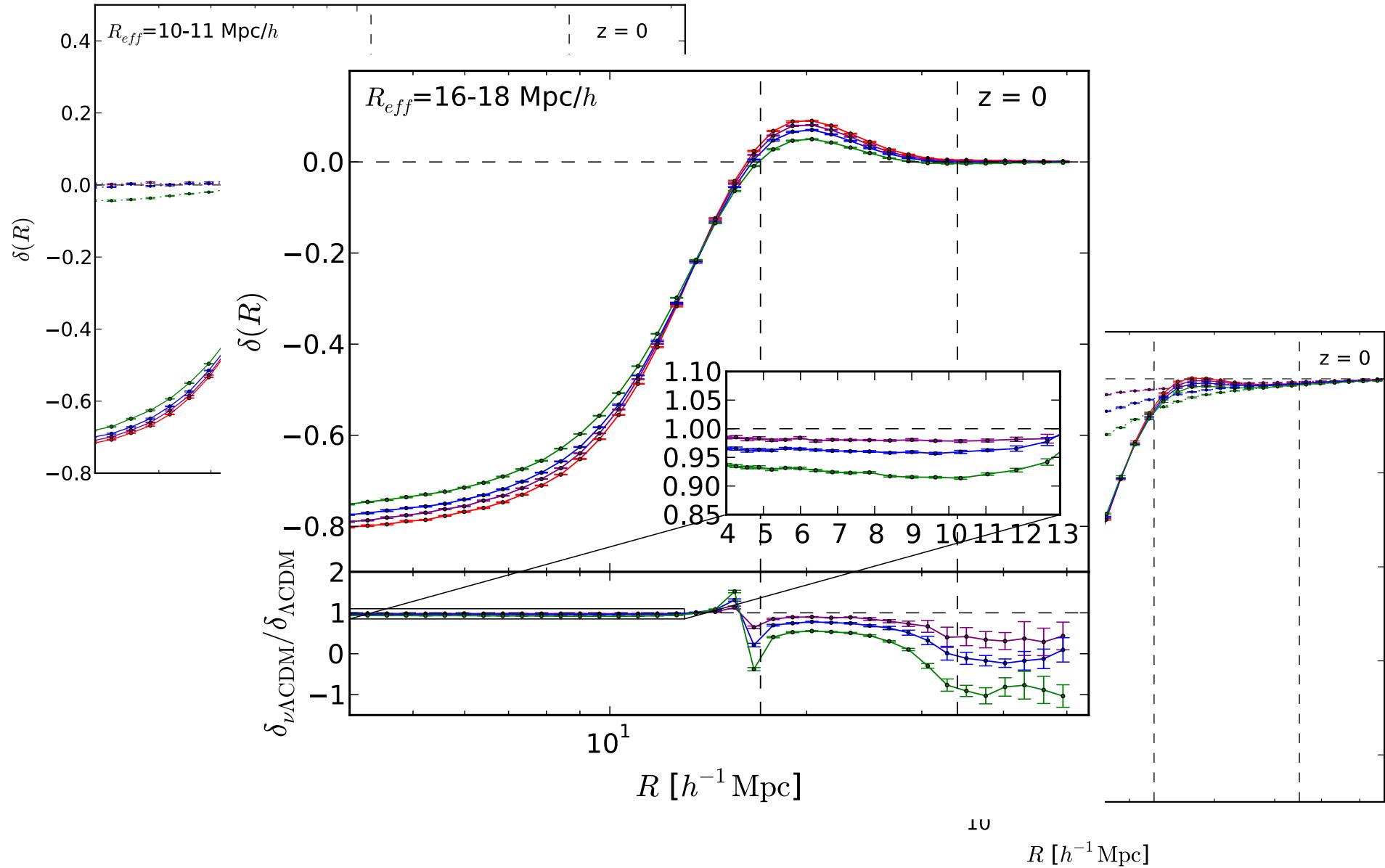
Neutrino effects on voids

Massara, FVN, Viel, Sutter 2015



Neutrino effects on voids

Massara, FVN, Viel, Sutter 2015



Future constraints from cosmology

Particle physics

$$m(\bar{\nu}_e) \leq 2.3 \text{ eV (95% C.L.)}$$

Kraus et al. 2005

Tritium beta decay

$$m(\bar{\nu}_e) \leq 0.2 \text{ eV (90% C.L.)}$$

KATRIN

Tritium beta decay

Cosmology

$$\sum_i m_{\nu_i} < 0.12 \text{ eV (95% C.L.)}$$

Palanque-Delabrouille et al. 2015

CMB + BAO + Lya forest

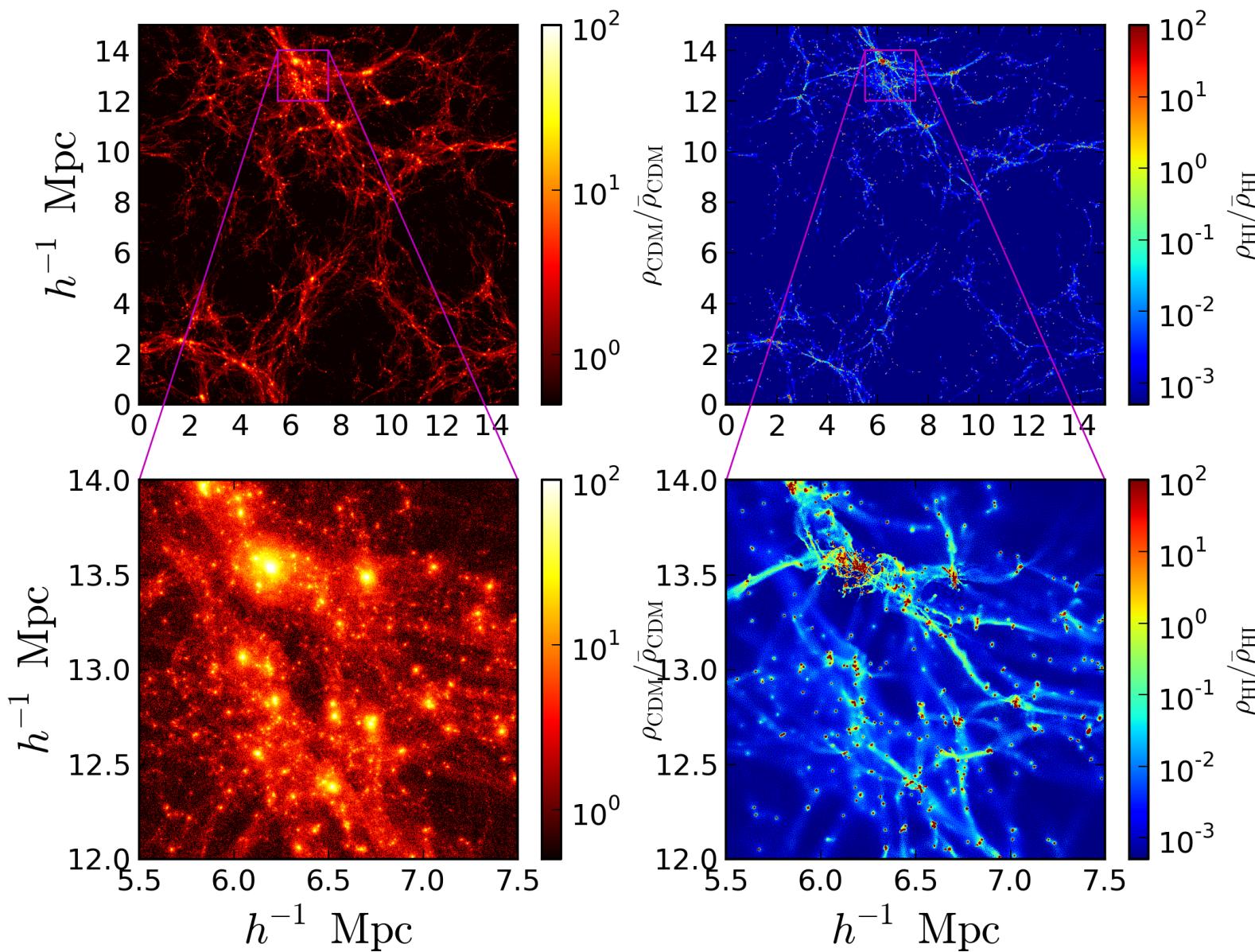
?

21cm

CMB+21cm

Weighing neutrinos with cosmic neutral hydrogen

Post-reionization era



Weighing neutrinos with cosmic neutral hydrogen

1. Study the properties of HI in cosmologies with massive neutrinos using hydrodynamic simulations.
2. Forecast the sensitivity of SKA to the sum of the neutrino masses.

Neutrino effects on 21cm

Hydrodynamical simulations

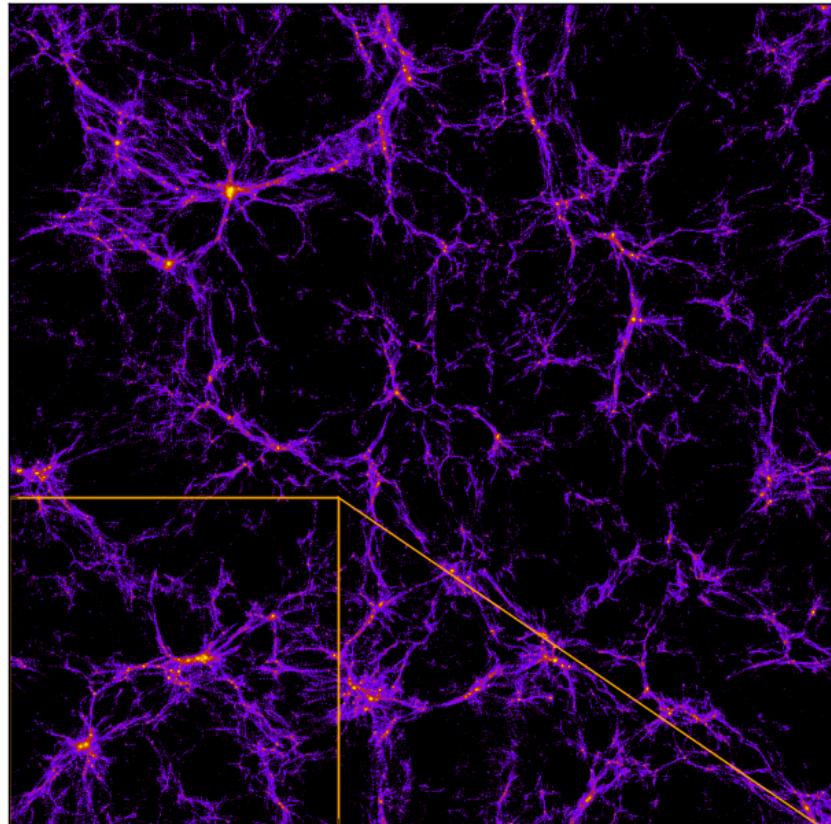
HI self-shielding

Presence of H₂

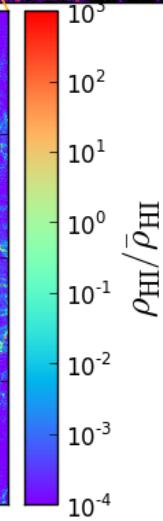
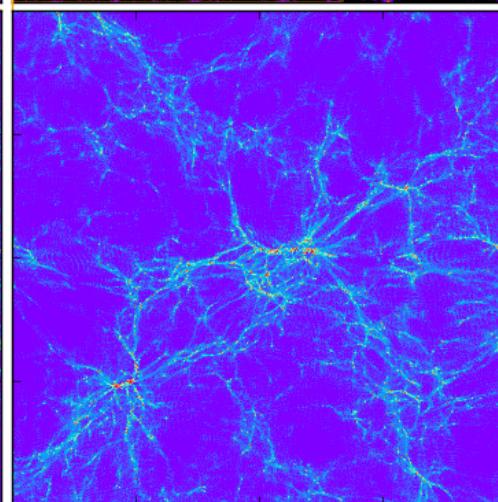
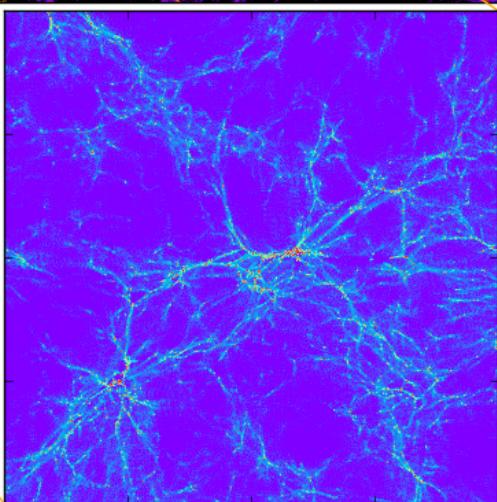
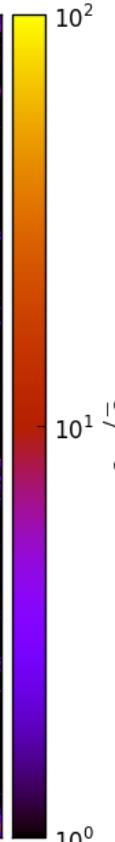
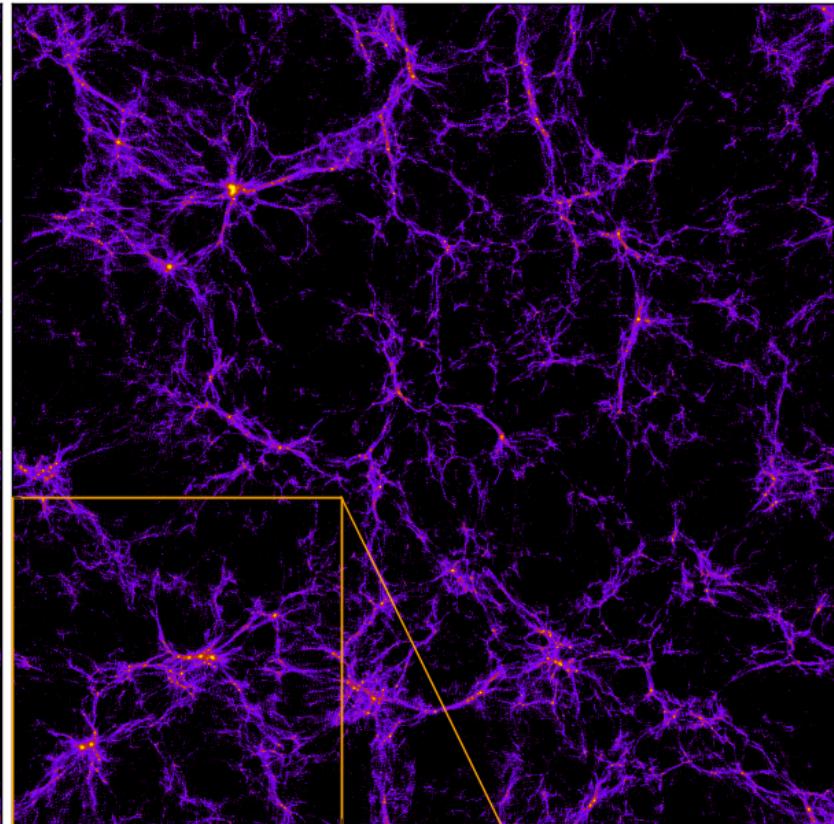
- Radiative cooling by H and He
- Heating by uniform UV background
- Star formation
- Feedback (galactic winds)
- CDM + baryons + neutrinos

Name	Box (h^{-1} Mpc)	Ω_{cdm}	Ω_b	Ω_ν	Ω_Λ	Ω_k	h	n_s	$10^9 A_s$	σ_8 ($z = 0$)
\mathcal{F}	50	0.2685	0.049	0.0	0.6825	0	0.67	0.9624	2.13	0.834
ν^+	50	0.2685	0.049	0.007075	0.675425	0	0.67	0.9624	2.13	0.778
ν_m^+	50	0.261425	0.049	0.007075	0.6825	0	0.67	0.9624	2.13	0.764
ν_m^{++}	50	0.25435	0.049	0.01415	0.6825	0	0.67	0.9624	2.13	0.693
\mathcal{C}^+	50	0.287	0.049	0.0	0.664	0	0.67	0.9624	2.13	0.868
\mathcal{C}^-	50	0.25	0.049	0.0	0.701	0	0.67	0.9624	2.13	0.797
\mathcal{B}^+	50	0.2685	0.055	0.0	0.6765	0	0.67	0.9624	2.13	0.816
\mathcal{B}^-	50	0.2685	0.043	0.0	0.6885	0	0.67	0.9624	2.13	0.853
\mathcal{H}^+	50	0.2685	0.049	0.0	0.6825	0	0.71	0.9624	2.13	0.886
\mathcal{H}^-	50	0.2685	0.049	0.0	0.6825	0	0.63	0.9624	2.13	0.777
\mathcal{N}^+	50	0.2685	0.049	0.0	0.6825	0	0.67	1.0009	2.13	0.846
\mathcal{N}^-	50	0.2685	0.049	0.0	0.6825	0	0.67	0.9239	2.13	0.822
\mathcal{A}^+	50	0.2685	0.049	0.0	0.6825	0	0.67	0.9624	2.45	0.894
\mathcal{A}^-	50	0.2685	0.049	0.0	0.6825	0	0.67	0.9624	1.81	0.769

$$\sum m_\nu = 0.0 \text{ eV}$$



$$\sum m_\nu = 0.6 \text{ eV}$$



Ingredients for 21cm IM

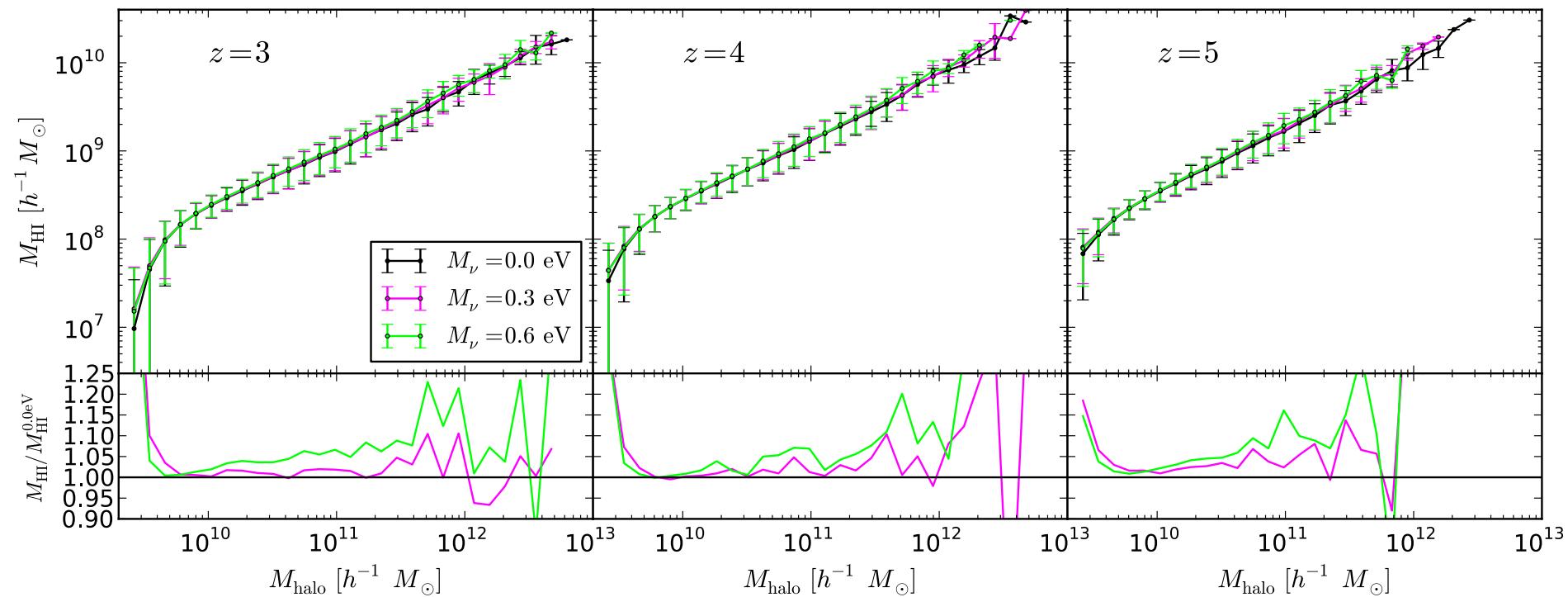
$$M_{HI}(M, z) \left\{ \begin{array}{l} b_{HI}(z) = \frac{\int\limits_0^{\infty} n(M, z) b(M, z) M_{HI}(M, z) dM}{\int\limits_0^{\infty} n(M, z) M_{HI}(M, z) dM} \\ \Omega_{HI}(z) = \frac{1}{\rho_c} \int\limits_0^{\infty} n(M, z) M_{HI}(M, z) dM \end{array} \right.$$

$$\overline{\delta T_b}(z) = 189 \left(\frac{H_0 (1+z)^2}{H(z)} \right) \Omega_{HI}(z) h \text{ mK}$$

$$P_{21cm}(k, z) = \overline{\delta T_b^2}(z) b_{HI}^2(z) \left(1 + \frac{2}{3} \beta(z) + \frac{1}{5} \beta^2(z) \right) P_m(k, z)$$

Neutrino effects on 21cm

FVN, Bull, Viel 2015



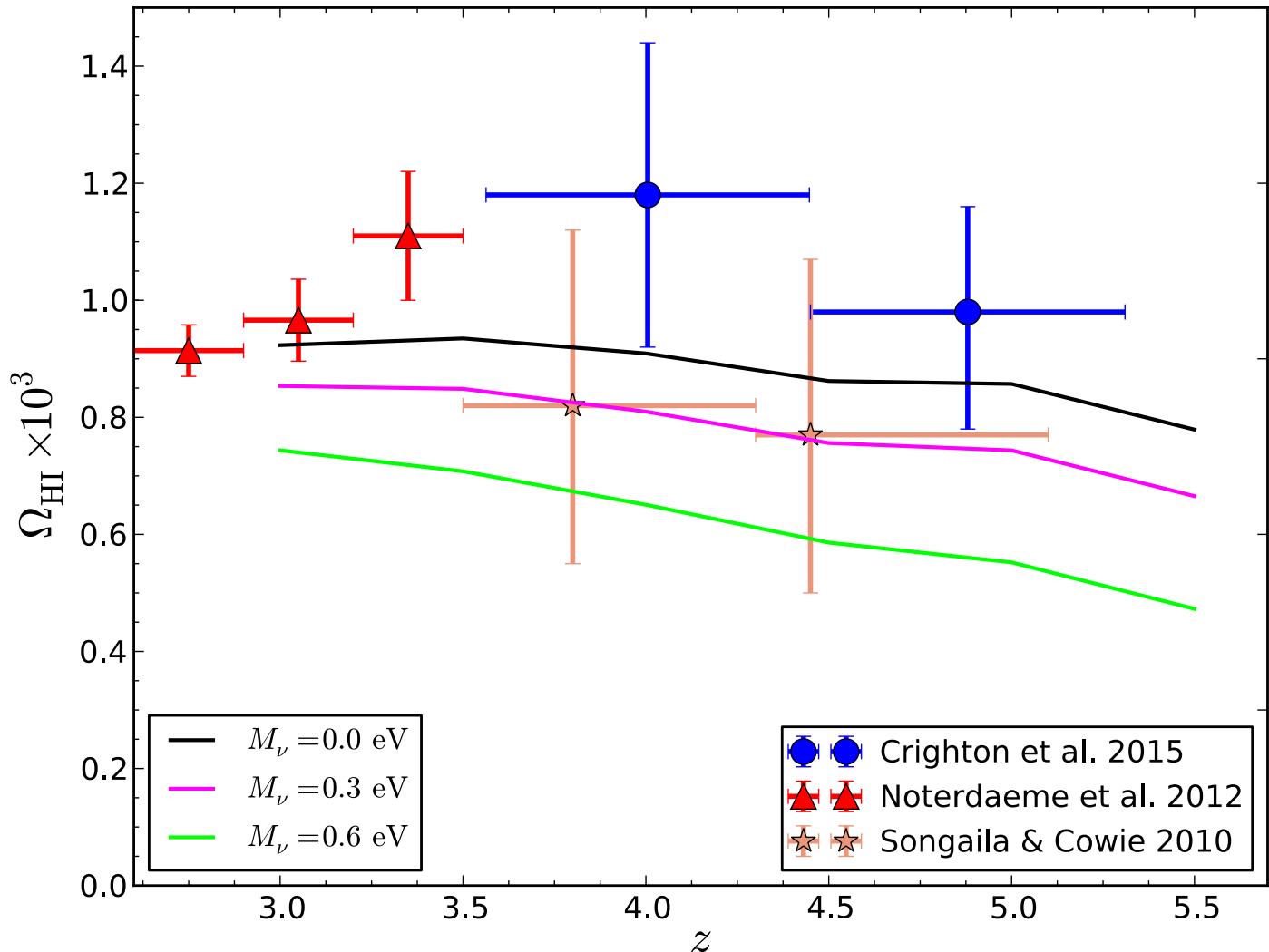
Halos with low masses do not host HI.

$M_{\text{HI}}(M, z)$ will exhibit a cut-off at low masses

Neutrino effects on 21cm

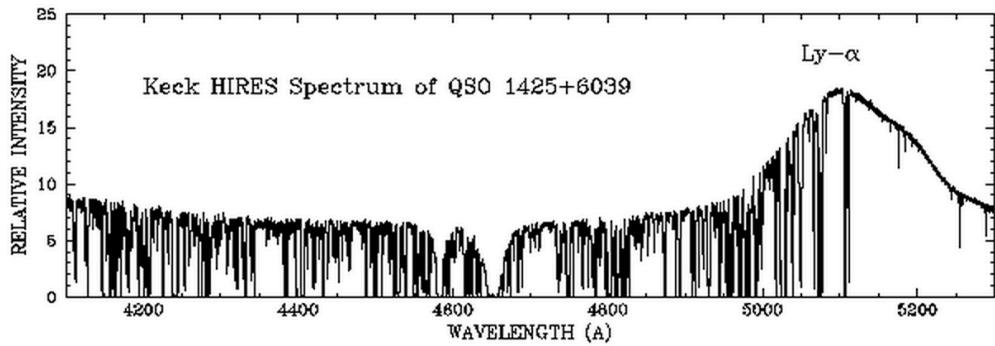
$$\Omega_{HI}(z) = \frac{1}{\rho_c^0} \int_0^\infty n(M, z) M_{HI}(M, z) dM$$

FVN, Bull, Viel 2015

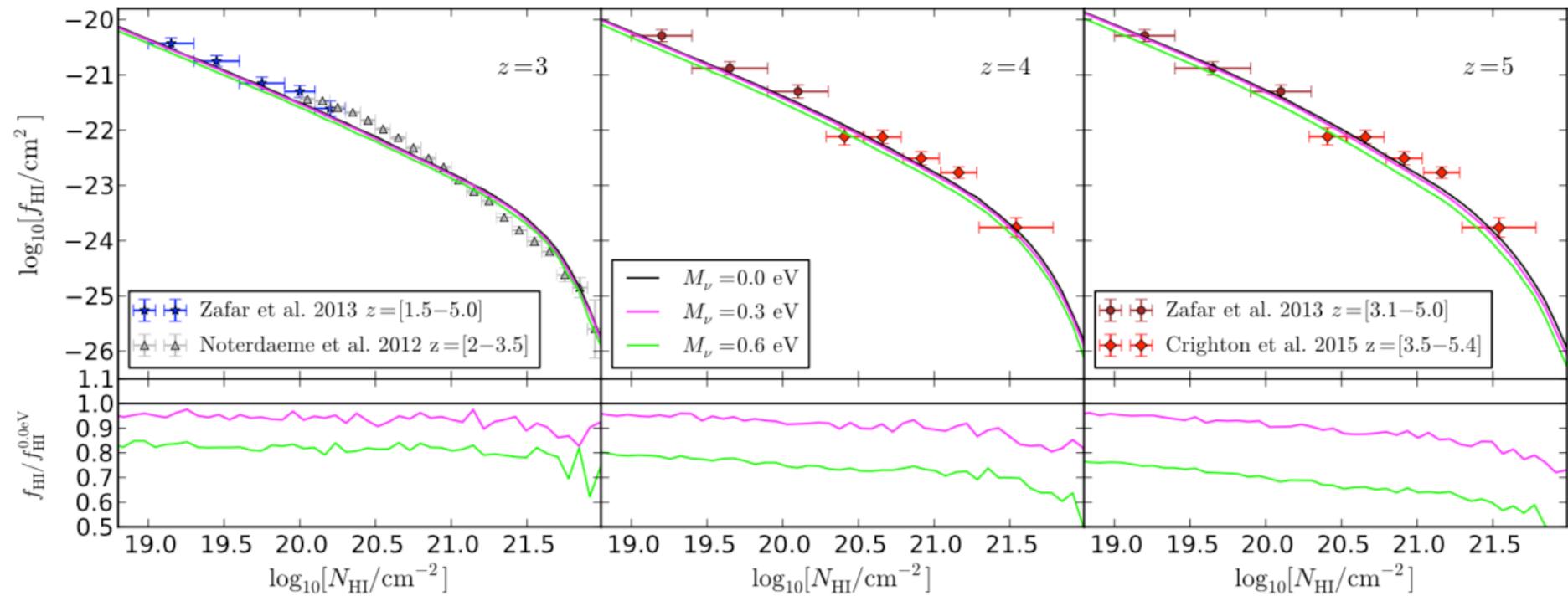


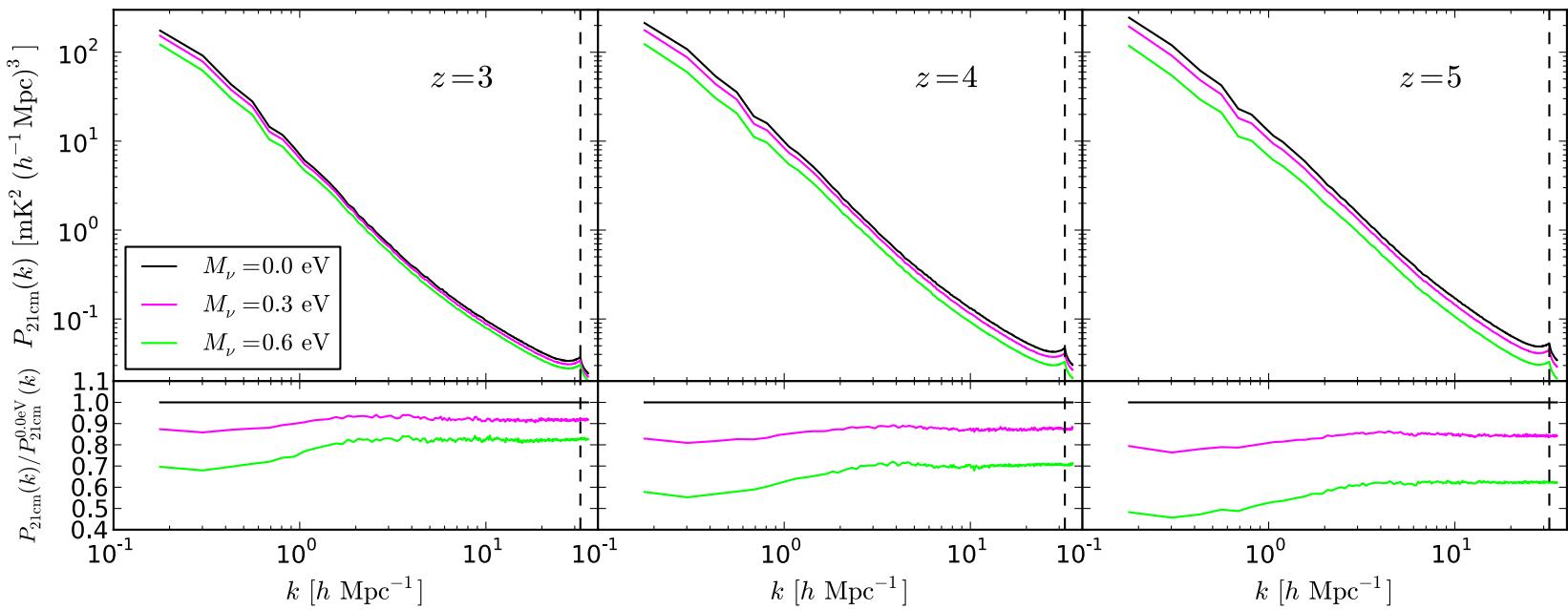
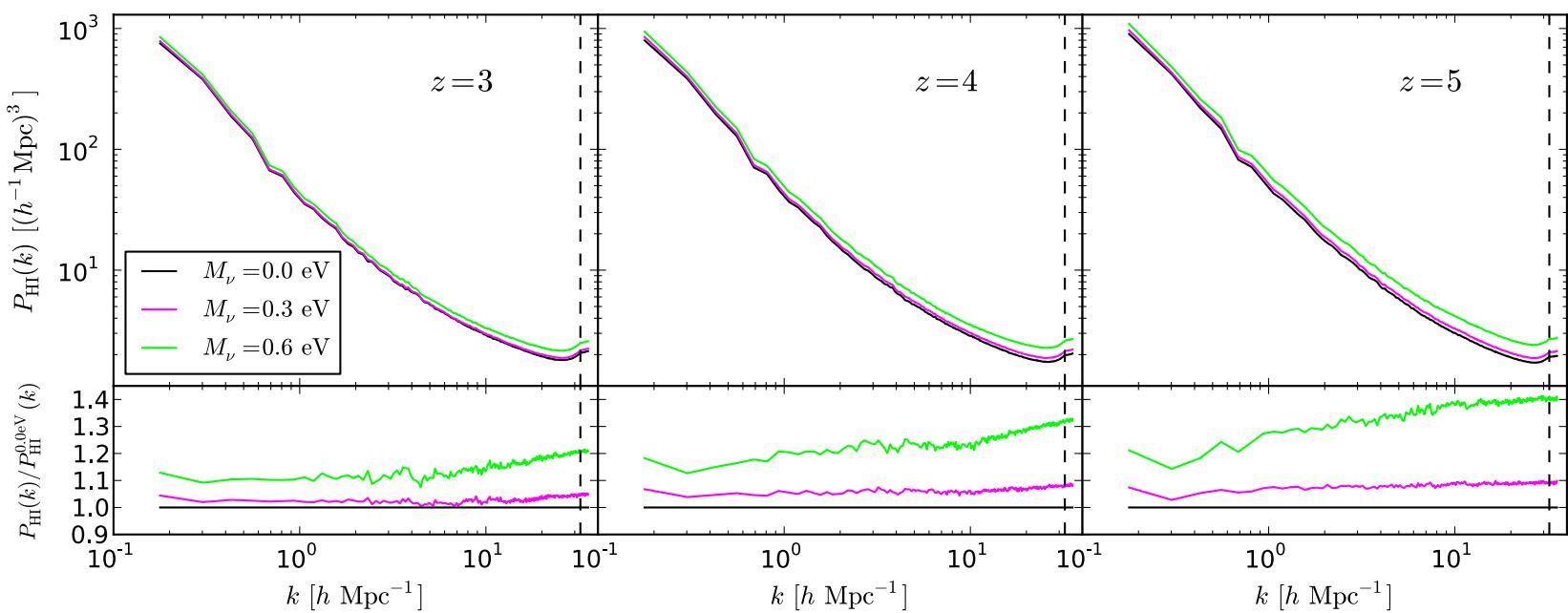
Neutrino effects on 21cm

FVN, Bull, Viel 2015



$$f_{HI}(N_{HI}) = \frac{d^2n}{dN_{HI}dX}$$

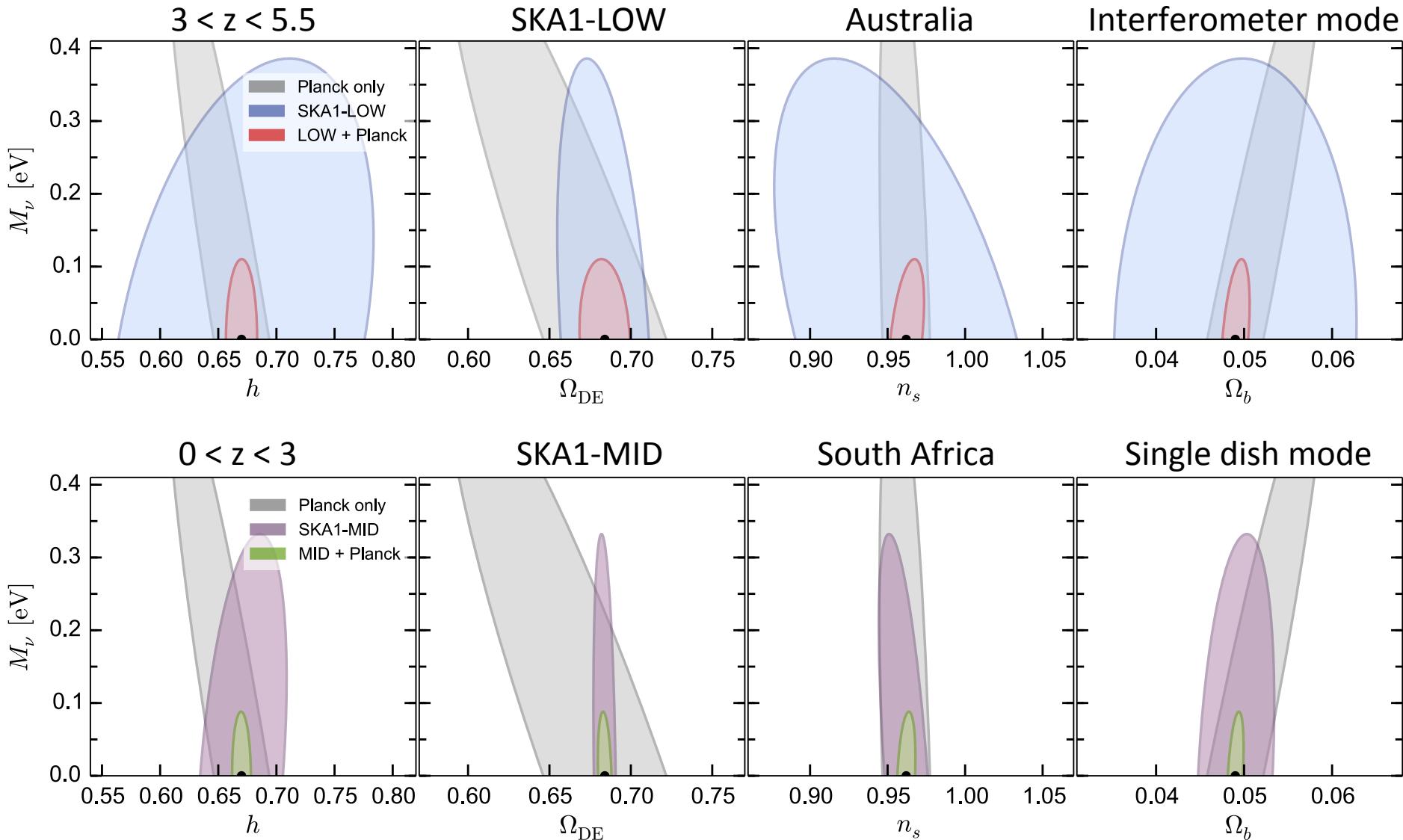




$$P_{21cm}(k, z) = \overline{\delta T_b^2}(z) b_{HI}^2(z) \left(1 + \frac{2}{3} \beta(z) + \frac{1}{5} \beta^2(z) \right) P_m(k, z)$$

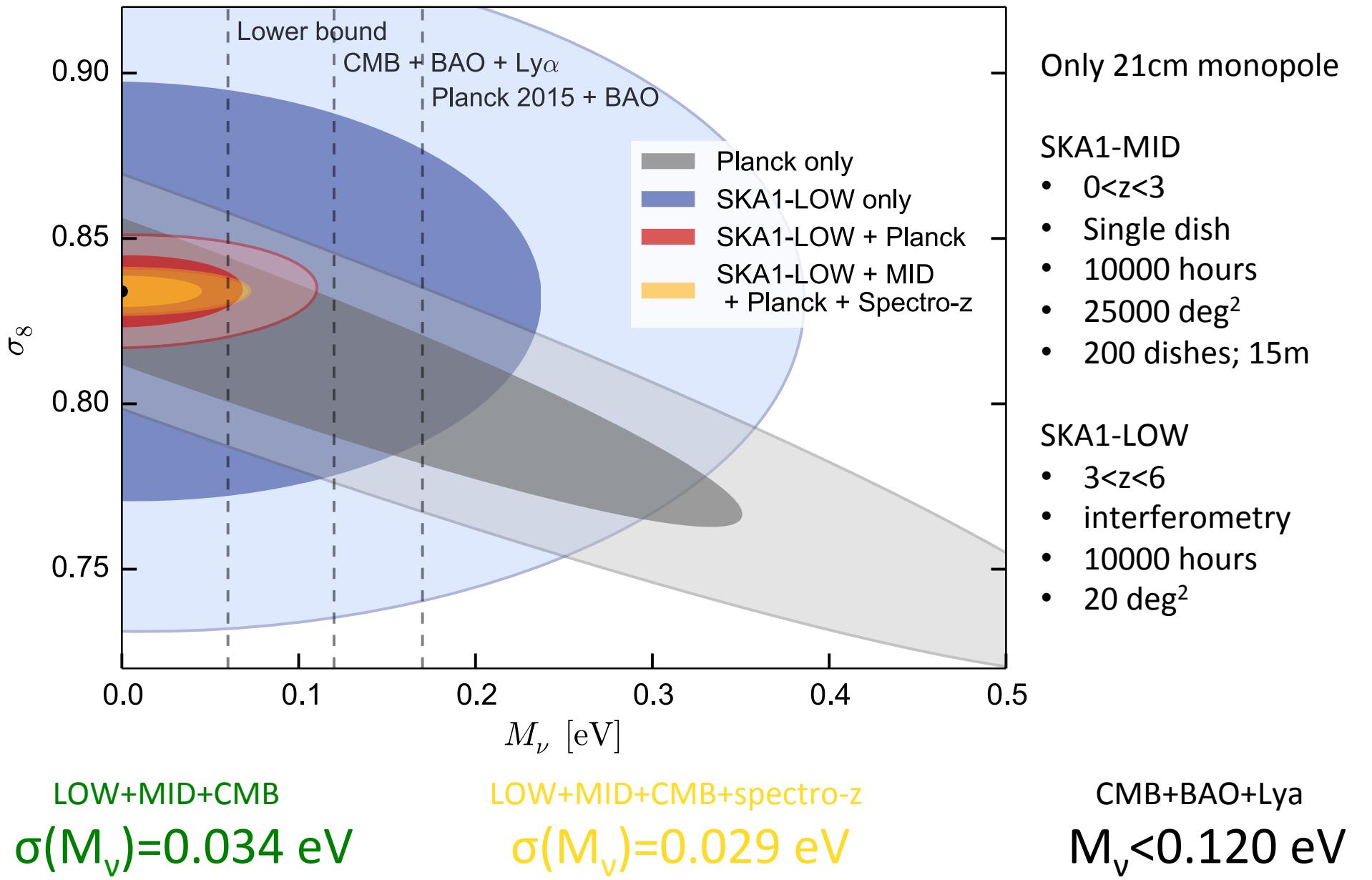
Forecasts for SKA

FVN, Bull, Viel 2015



Forecasts for SKA

FVN, Bull, Viel 2015



Conclusions

- Neutrinos have mass!!!
 - Major consequences for particle physics and cosmology
- Big Bang theory predicts the existence of the CvB
- Massive neutrino effects well understood at linear order
- Many effects at fully non-linear level
 - Neutrino clustering
 - Halo properties
 - Halo mass function
 - Bias
 - BAO
 - Voids
 - HI abundance and clustering
- SKA IM survey+CMB will set $\sigma(M_\nu) = 0.034 \text{ eV}$ ($M_\nu < 0.120 \text{ eV}$)